

3rd WORKSHOP ON PARTICLE PHYSICS

NATIONAL CENTRE FOR PHYSICS
(QUAID-I-AZAM UNIVERSITY)

Detectors for High Energy Physics

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<http://rolandi.home.cern.ch/rolandi/>

Literature on particle detectors (1)

Text books

- C. Grupen, Particle Detectors, Cambridge University Press, 1996
- G. Knoll, Radiation Detection and Measurement, 3rd Edition, 2000
- W. R. Leo, Techniques for Nuclear and Particle Physics Experiments, 2nd edition, Springer, 1994
- R.S. Gilmore, Single particle detection and measurement, Taylor&Francis, 1992
- W. Blum, G. Rolandi, Particle Detection with Drift Chambers, Springer, 1994
- K. Kleinknecht, Detektoren für Teilchenstrahlung, 3rd edition, Teubner, 1992

Literature on particle detectors (2)

Review articles

- **Experimental techniques in high energy physics**, T. Ferbel (editor), World Scientific, 1991.
- **Instrumentation in High Energy Physics**, F. Sauli (editor), World Scientific, 1992.
- Many excellent articles can be found in **Ann. Rev. Nucl. Part. Sci.**

Other sources

- Particle Data Book (Phys. Rev. D, Vol. 54, 1996)
- **R. Bock, A. Vasilescu**, Particle Data Briefbook
<http://www.cern.ch/Physics/ParticleDetector/BriefBook/>
- Proceedings of detector conferences (Vienna VCI, Elba, IEEE)

High Energy Physics Experiments

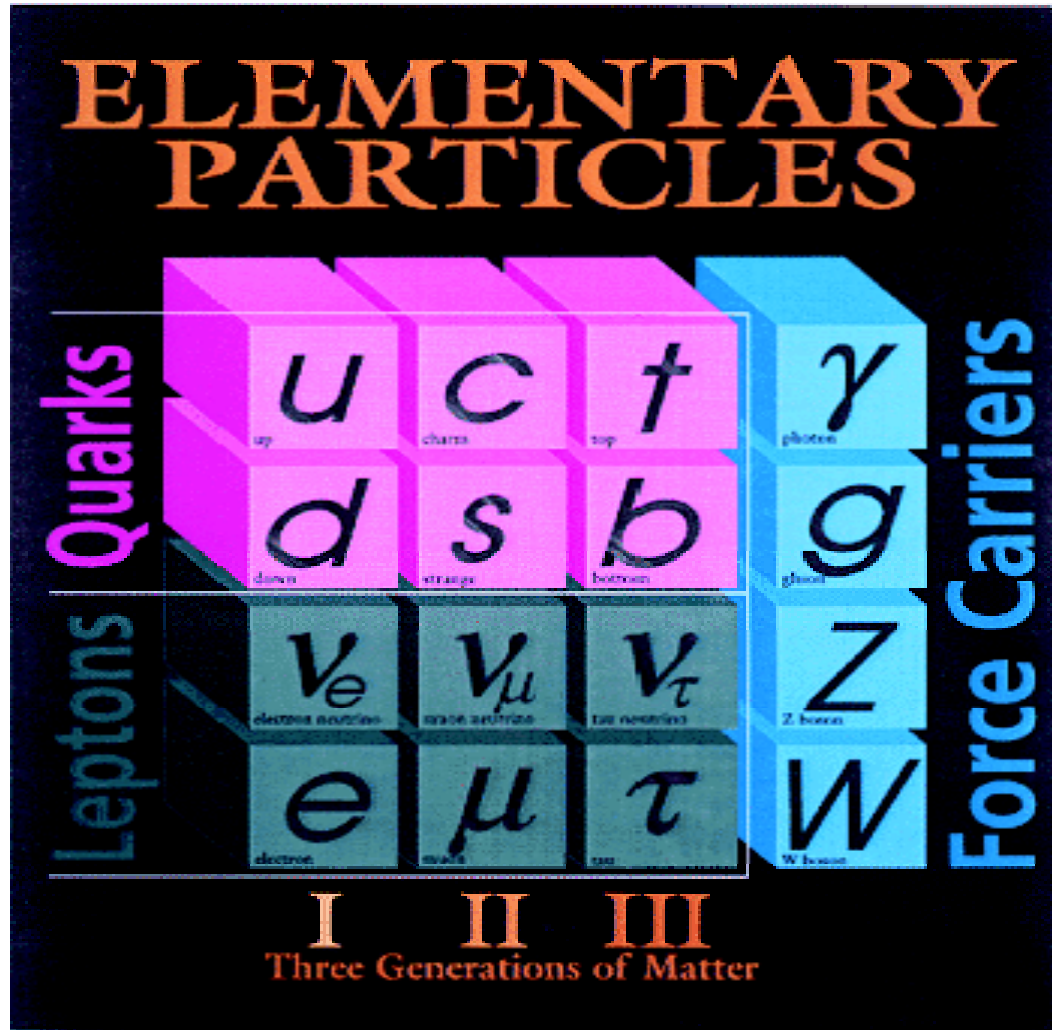
High Energy Physics studies the interactions between elementary particles:

QUARKS AND LEPTONS

These interactions are mediated by

FORCE CARRIERS

Elementary particles



Jets of particles

photons

Electrons,
Muons,
Missing E_T

Fundamental particles and interactions

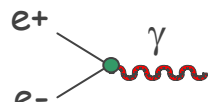
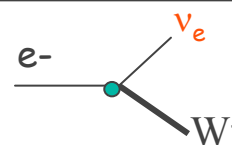
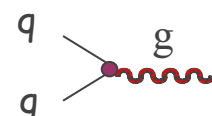
Matter particles : fermions, spin = 1/2

e	μ	τ	$q = -1$	u	c	t	$q = +2/3$
ν_e	ν_μ	ν_τ	$q = 0$	d	s	b	$q = -1/3$

+ anti-particles

Why 3 families ?
 Why fermion masses ?
 Why boson masses ?
 Are quarks elementary ?
 Why gauge symmetries ?
 etc ...

Interactions specified by symmetry : $U(1)_Y \times SU(2)_W \times SU(3)_C$

Force carriers : bosons, spin=1					
Particle	Force		Coupling (E~100 GeV)	Mass	Intensity
γ	EM (charged particles)		$\alpha_{EM} = \frac{e^2}{4\pi} \approx 0.008$	0	$\sim 10^{-1}$
W^\pm, Z	weak (q, l, W^\pm, Z)		$\alpha_W = \frac{g^2}{4\pi} \approx 0.03$	$\sim 100 \text{ GeV}$	$\sim 10^{-5}$
8 g	strong (q, g)		$\alpha_s = \frac{g_s^2}{4\pi} \approx 0.12$	0	1

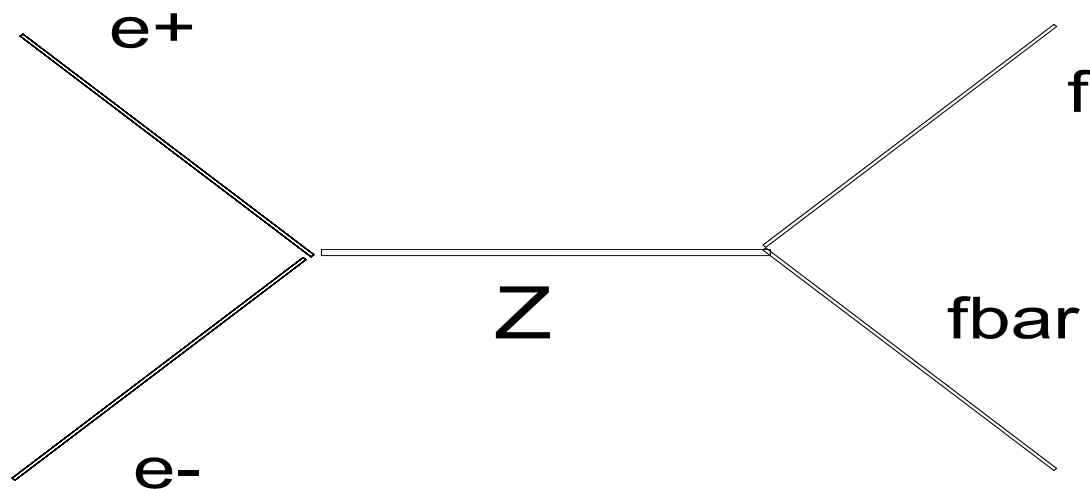
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to str

Mass "generator" : Higgs scalar, spin=0 ?
 (EWSB)
 predicted by SM but not yet observed

High Q² Reactions

High Q² reactions involve partons
(quarks and leptons)

Measure the properties
(parameters) of the partons in
the initial and final state

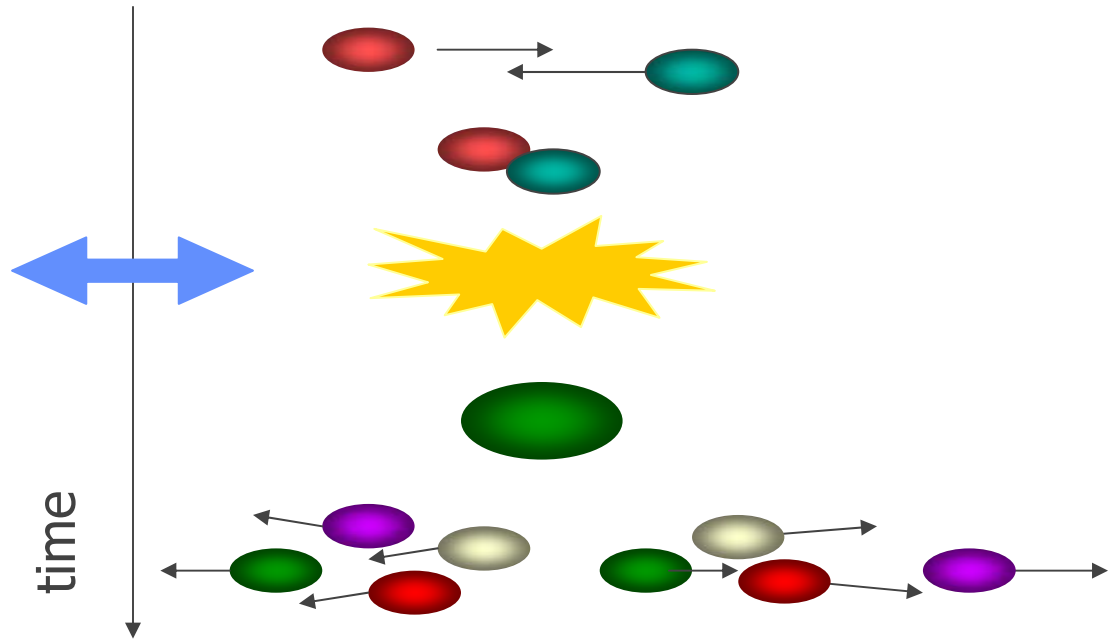
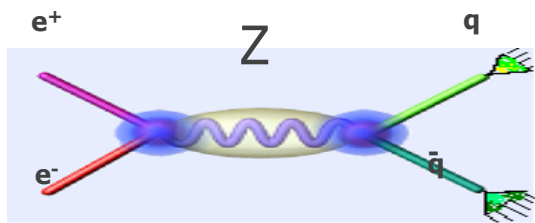


- Charge
- Direction
- Flavor
- Energy
- Spin

Realistic views of an elementary particle reaction

$$e^+ + e^- \rightarrow Z^0 \rightarrow q \bar{q}$$

Quarks hadronize in jets



Usually we can only 'see' the end products of the reaction, but not the reaction itself.

In order to reconstruct the reaction mechanism and the properties of the involved particles, we want the maximum information about the initial state and the end products !

Colliding beams

Colliding beams are the most efficient way to maximize the center of mass energy of the collision between the partons

In order to store and collide efficiently the beams, only charged and stable particle can be used:

Electrons and positrons

Protons and anti-protons

Nuclei



(muon colliders are also under study)

Luminosity

particles are grouped in packets called bunches

terminology:

- N_p = # protons in bunch
- $N_{\bar{p}}$ = # p-bar in bunch
- B = # bunches
- f = revolution freq.
-

L scales with E_{beam}

- Smaller beams at higher energies

Protons in Bunch

Total Antiprotons

Revolution Frequency

$$L = \frac{f N_p (B N_{\bar{p}})}{2\pi (\sigma_p^2 + \sigma_{\bar{p}}^2)} F(\sigma_z / \beta^*)$$

Beam Sizes

Beam Shape Form factor at Intersection

Luminosity

Luminosity, L , is a measurement of the brightness of the interaction region

The luminosity is a major machine parameter

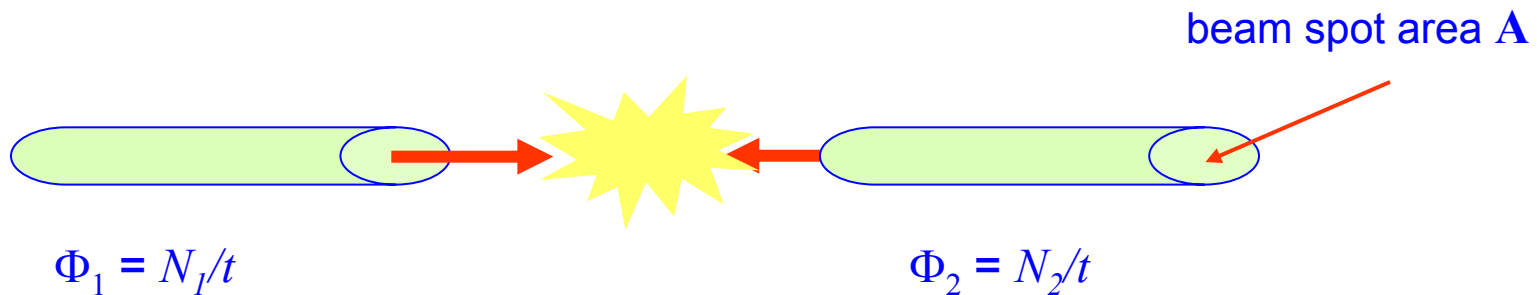
- High luminosity \rightarrow sensitivity to rare events



The concept of cross sections

Cross sections σ or differential cross sections $d\sigma/d\Omega$ are used to express the probability of interactions between elementary particles.

Example: 2 colliding particle beams



What is the interaction rate R_{int} ?

$$R_{int} \propto N_1 N_2 / (A \cdot t) = \sigma \cdot L$$

Luminosity L [$\text{cm}^{-2} \text{s}^{-1}$]

σ has dimension area !

Practical unit:

$$1 \text{ barn (b)} = 10^{-24} \text{ cm}^2$$

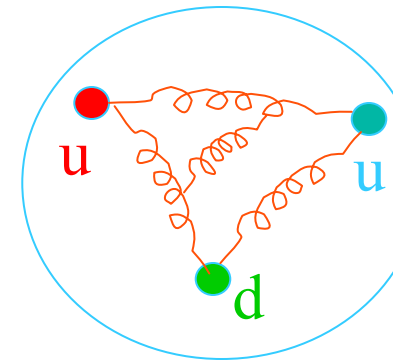
Different beams properties

electrons are point-like

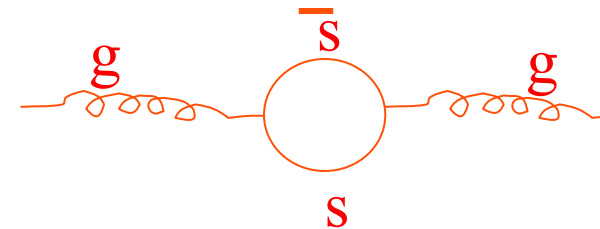
protons are a quark-gluon soup

- 3 valence quarks bound by exchange of gluons
 - Gluons are colored and interact with other gluons
 - Virtual quark pair loops can pop-up generating additional quark content (sea-quarks)
- Proton momentum is shared among all constituent partons (quarks & gluons)

effective parton collision energy is lower than available CM energy and depends on the proton structure described by pdf



Proton



Virtual quark loop

Proton structure

parton distribution functions (pdf) describe the momentum distribution of each parton species in the proton

– When protons collide their structure is affected

- Pdf's depend on Q of interaction

– $F_i(x, Q)$

- $x = p_{\text{parton}}/p_{\text{proton}} =$ fractional momentum carried by parton
- $Q =$ momentum transfer of interaction
- $i =$ parton species

Effective coll. Energy (Teva)

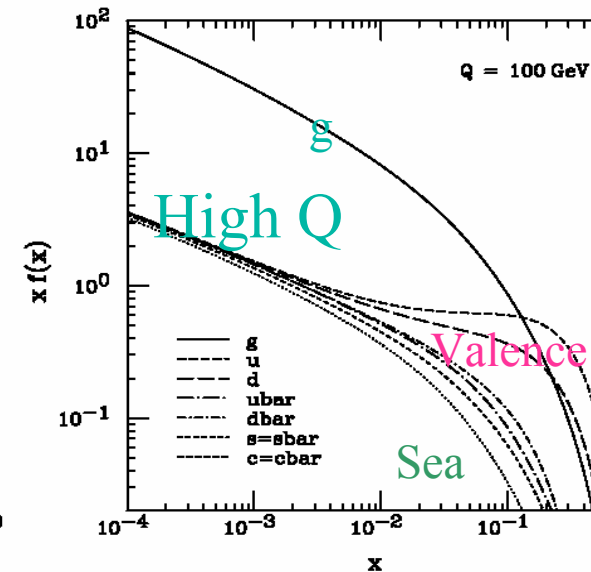
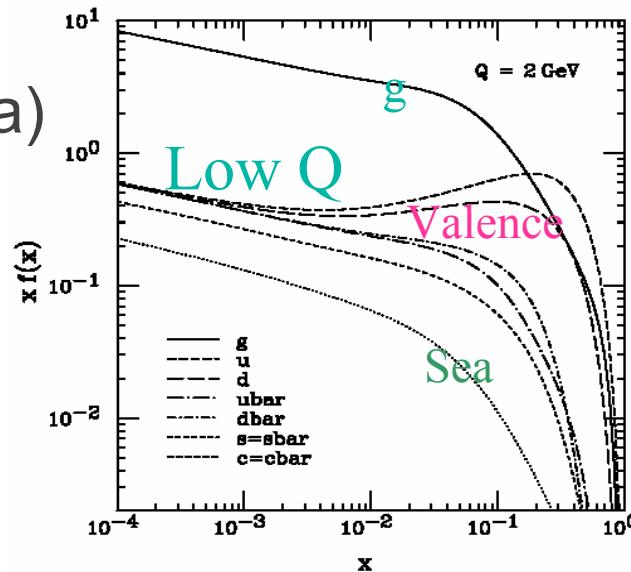
– $E_{\text{parton}}^{\text{CM}} = \sqrt{x_1 x_2} E^{\text{CM}}$

– $\langle x \rangle \sim 0.1$

– $\langle E_{\text{parton}}^{\text{CM}} \rangle \sim 200 \text{ GeV}$

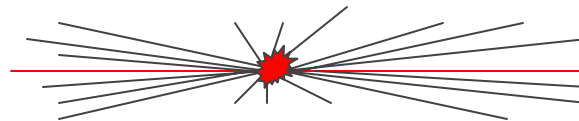
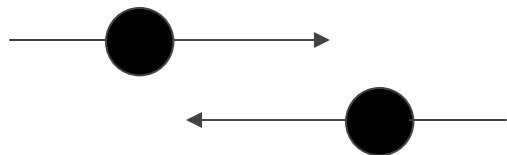
• **Max $E_{\text{parton}}^{\text{CM}}$**

Observed 1364 GeV



Phenomenology of pp collisions

Most interactions due to collisions at large distance between incoming protons where protons interact as “ a whole ” → small momentum transfer ($Dp \approx \eta / Dx$) → particles in final state have large longitudinal momentum but small transverse momentum (scattering at large angle is small)

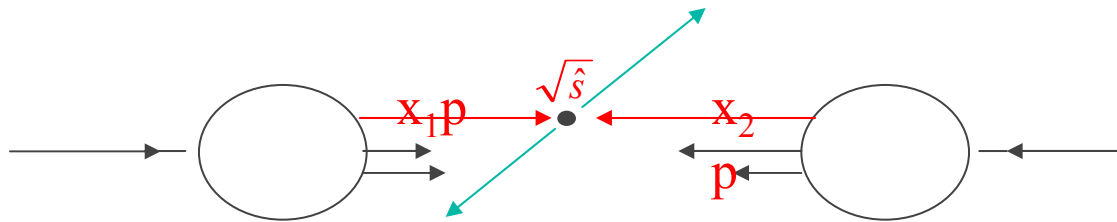


charged particles uniformly distributed in ϕ Most energy escapes down the beam pipe

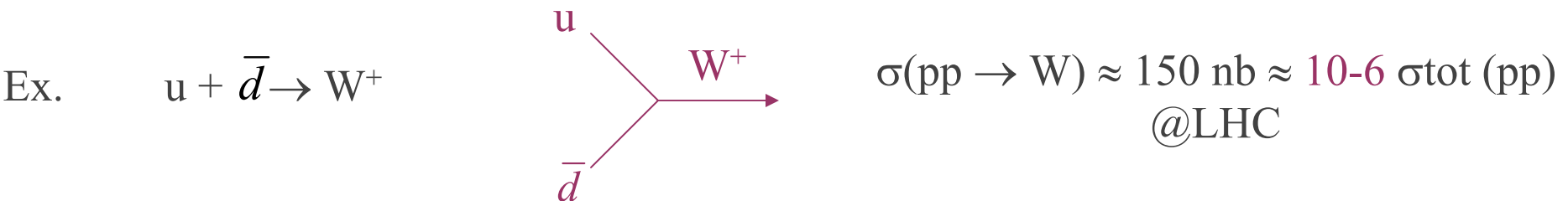
These are called minimum-bias events (“ soft “ events). They are the large majority but are not very interesting

Phenomenology of pp collisions

Monochromatic proton beam can be seen as **beam of quarks and gluons** with a wide band of energy. Occasionally **hard scattering** (“**head on**”) between constituents of incoming protons occurs.



Interactions at **small distance** → **large momentum transfer** → **massive particles** and/or **particles at large angle** are produced. These are **interesting physics events** but they are **rare**.



Hadron Colliders vs. e+e-

e+e- storage rings: **best for precision measurements**

- Very clean and well defined final state.... **However**
 - Max energy limited by electron radiation:
 - Energy loss: $8.85 \times 10^{-5} E^4/\rho$ MeV/turn (E is in GeV, ρ in km)
 - » **Scales with 4th power of particle mass**
 - Max LEP CM energy ~200 GeV
- **Only final states which couple to photon or Z boson at precise CM energy**

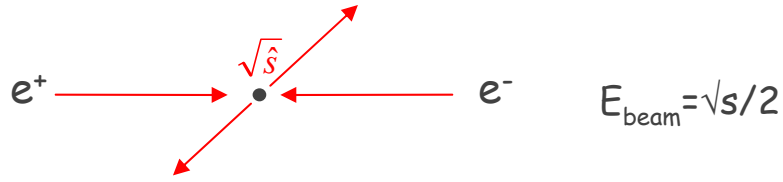
hadron colliders: **best as discovery machines**

- Final state is more complex..... **However**
 - Much higher energy available
 - Energy loss: $7.8 \times 10^{-18} E^4/\rho$ MeV/turn (E is in GeV, ρ in km)
 - Max **Tevatron CM energy ~2,000 GeV** (14,000 GeV at LHC !)
- **Broad band collisions of many different initial partons**

e^+e^- Colliders

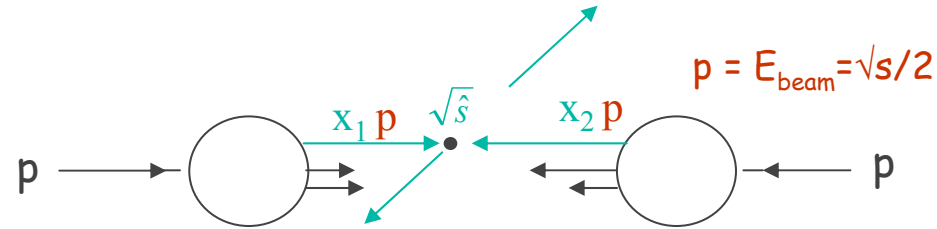
vs

pp/pp Colliders



Energy of elementary interaction known

$$\sqrt{\hat{s}} = E(e^-) + E(e^+) = \sqrt{s}$$



• Energy of elementary interaction not known

$$\sqrt{\hat{s}} = \sqrt{x_1 x_2 s} < \sqrt{s}$$

Only two elementary particles collide
→ **clean final states**

• Elementary interaction (hard) + interaction of "spectator" q, g (soft) overlapped in detector

Mainly EW processes

• EW processes suffer from **huge backgrounds from strong processes**

\sqrt{s} limited by e^\pm synchrotron radiation:

$$E_{\text{loss}} \sim \frac{E_{\text{beam}}^4}{R} \frac{1}{m_e^4} \quad E_{\text{loss}} \sim 2.5 \text{ GeV/turn}$$

LEP2 ($E_{\text{beam}} \sim 100 \text{ GeV}$)

• Synchrotron radiation is $\sim (m_p/m_e)^4 \sim 10^{13}$ smaller

↓
high energy more difficult
→ next machine : Linear Collider
(TESLA, NLC, JLC, $\sqrt{s} = 500-800 \text{ GeV}$?)

clean environment → **precision measurements machines**

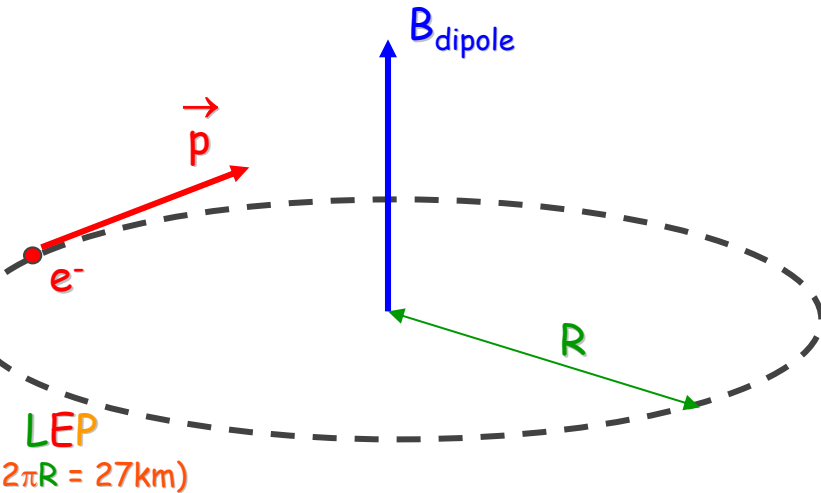
↓
-- **high energy easier** → **discovery machines**
next machine : LHC, pp, $\sqrt{s} = 14 \text{ TeV}$
in the LEP ring
-- "dirty" environment

Parameters of some colliders

	PEP II	LEP	Hera	Tevatron	LHC
	e ⁺ e ⁻	e ⁺ e ⁻	e p	p pbar	p p
Circumference (km)	2.2	26.6	6.3	6.3	26.6
Peak magnetic field (T)	0.18 and 0.75	0.135	0.27 and 4.6	4.4	8.3
Number of Dipoles	192	3280	396 and 413	774	1232
Maximum beam energy (GeV)	4 and 12	100	30 and 920	1000	7000
Luminosity (10 ³⁰ cm ⁻² s ⁻¹)	4600	24	75	210	10000
Time between collisions (μs)	0.0042	22	0.096	0.396	0.025
Energy spread (units 10 ⁻³)	0.7	1	1 and 0.2	0.09	0.1
Bunch length (cm)	1.1	1	1 and 8.5	38	7.5
Beam Radius (μm)	460 x 4	250 x 4	270 x 50	34 and 29	16
Particles per bunch (units 10 ¹⁰)	2.1 and 5.9	45	3 and 7	27 and 7.5	11
Average current (μAmp)	800 and 1100	5	40 and 90	81 and 22	536
Beam Polarization (%)	none	55	50	none	none

Measurement of the LEP Beam Energy (I)

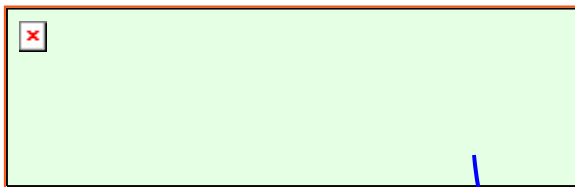
Approximation: LEP is a circular ring immersed in a uniform magnetic field:



$$E \sim p = e B R = (e/2\pi) B L$$

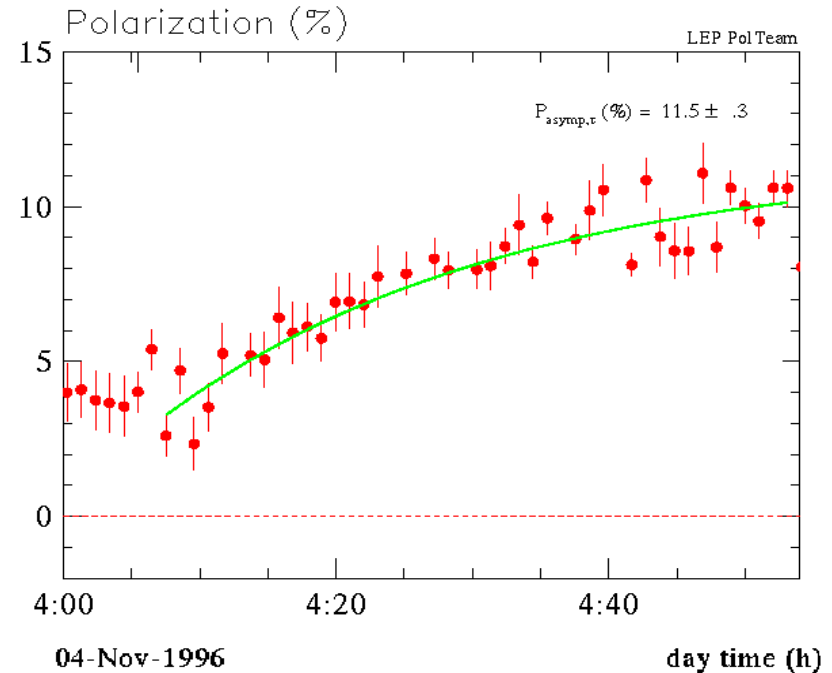
In real life:

B non-uniform, ring not circular



To be measured

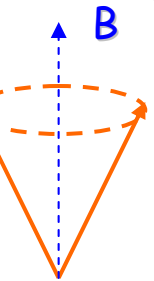
- 1) The electrons get **transversally polarized** (i.e., their spin tends to align with B), but
 - Process **very sensitive** to imperfections (→ slow, typically hours, and limited to o(10%) polarization)



- Process **very sensitive** to beam-beam interactions (→ one beam, no polarization in collisions)

Measurement of the LEP Beam Energy (II)

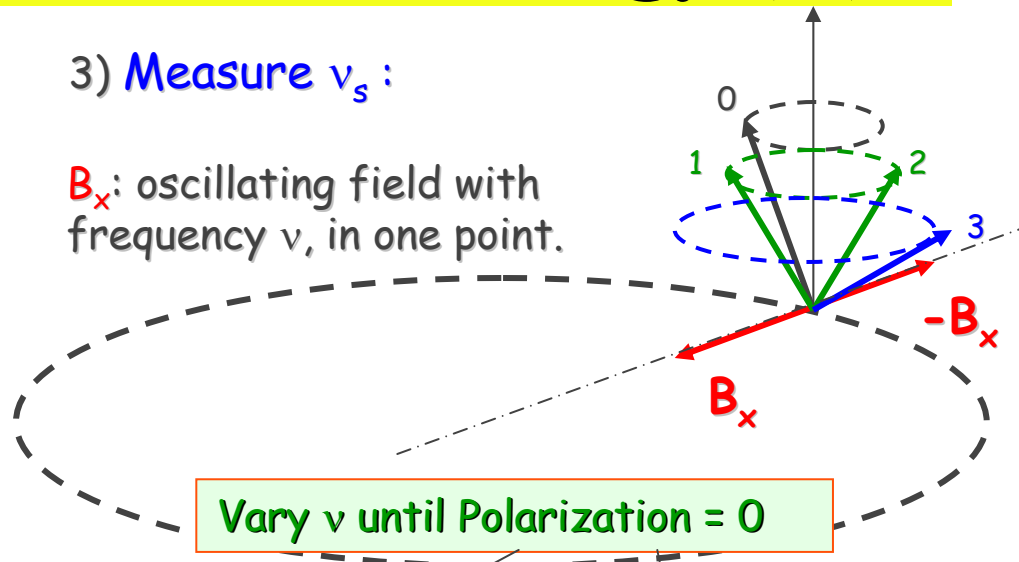
2) The spin precesses around B with a frequency proportional to B .



The number of revolutions for each LEP turn is thus proportional to BL (in fact, to $\int B dl$, and then to E_{beam})

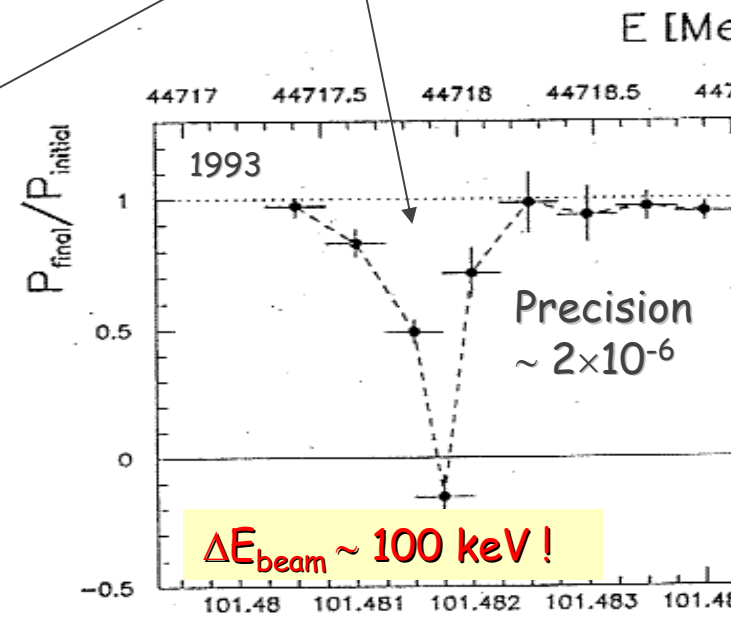
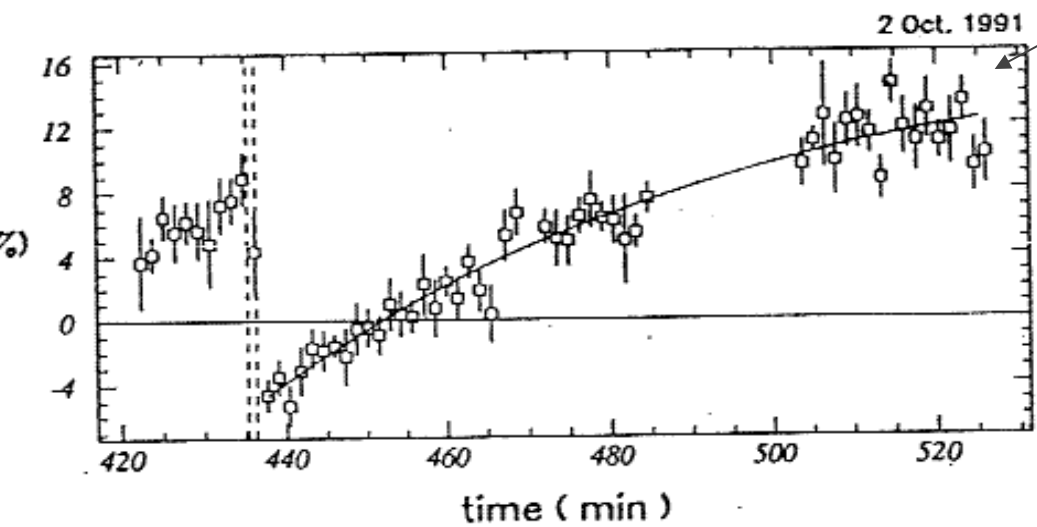
3) Measure ν_s :

B_x : oscillating field with frequency ν , in one point.



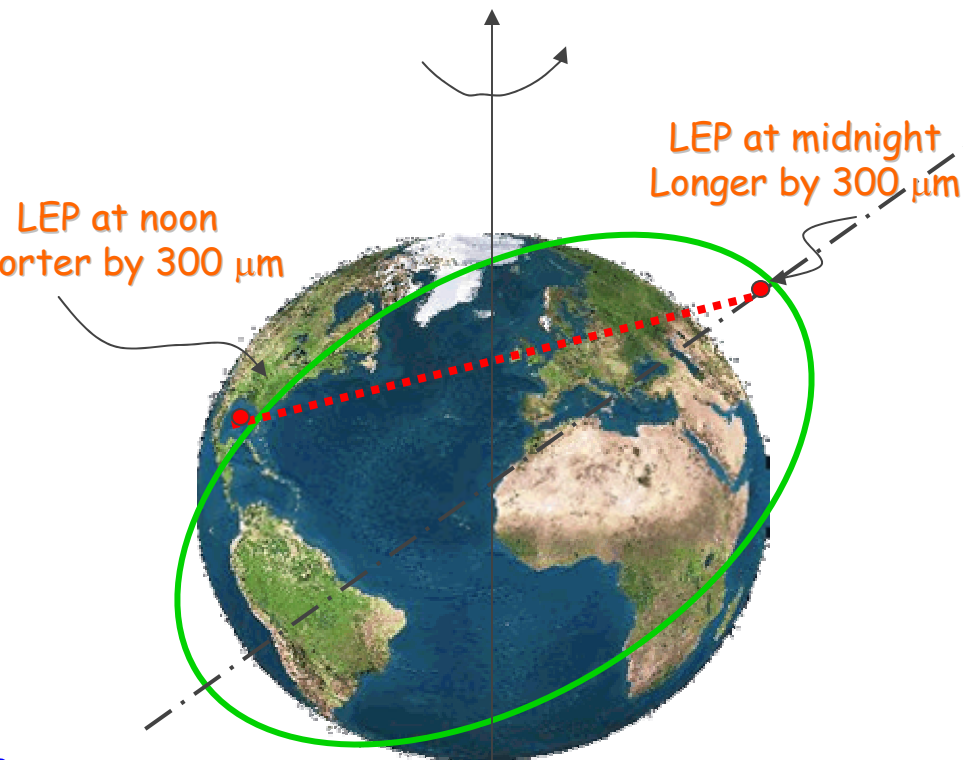
$$\nu_s = \frac{g_e - 2}{2m_e} \times E_{beam}$$

101.5	Peak-2
103.5	Peak
105.5	Peak+2



Measurement of the LEP Beam Energy (III)

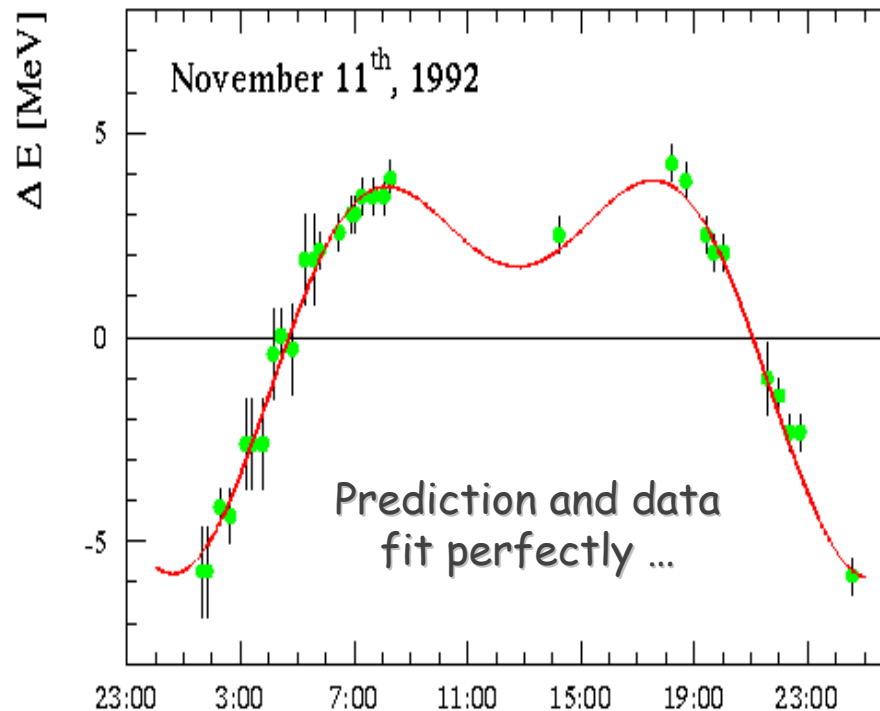
A dispersion of 10 MeV is observed ($\gg 100$ keV) at the same machine conditions. Correlation with the moon found on 1992, Nov 11th:



- At midnight, the electrons see **less** magnetic field, E is **smaller**;
- At noon, they see **more** magnetic field, and E is **larger**.

However, the electron orbit length is fixed by the RF frequency:

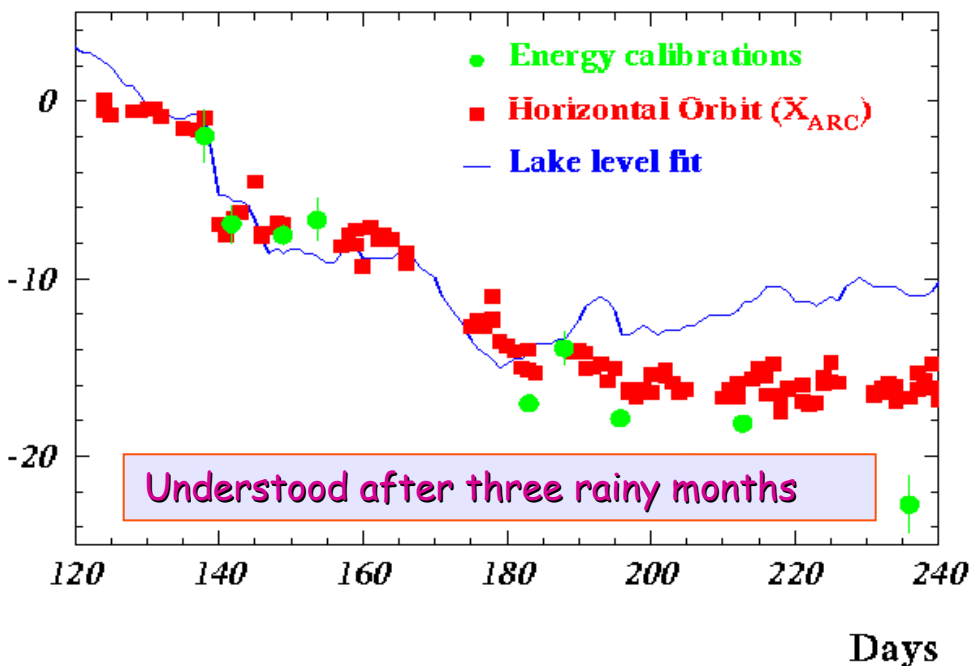
$$L = c \times \Delta t$$



Measurement of the LEP Beam Energy (IV)

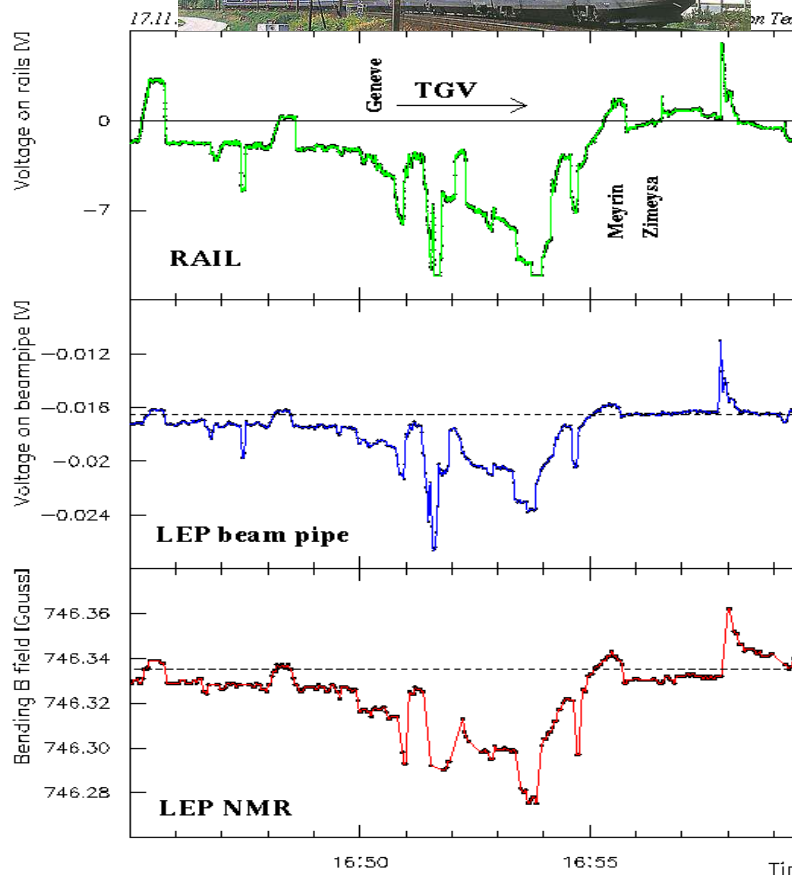
Other 10 MeV-ish effects understood even later:

Geological deformation due to the level of the lake
 rain change LEP circumference;
 (controlled with the BOM's)



Effect of the TGV: currents induced on the
 LEP beam pipe induce changes in the magnetic field
 (controlled by 16 NMR probes)

Understood after one-day strike

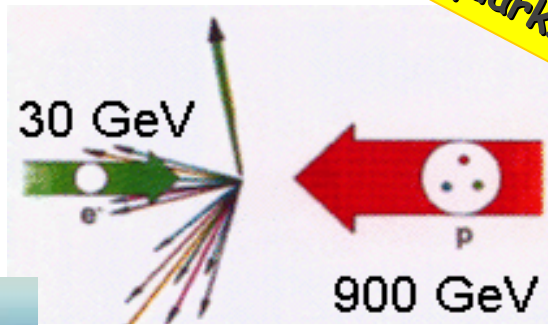


Now: $\Delta E_{beam} < 2 \text{ MeV}$

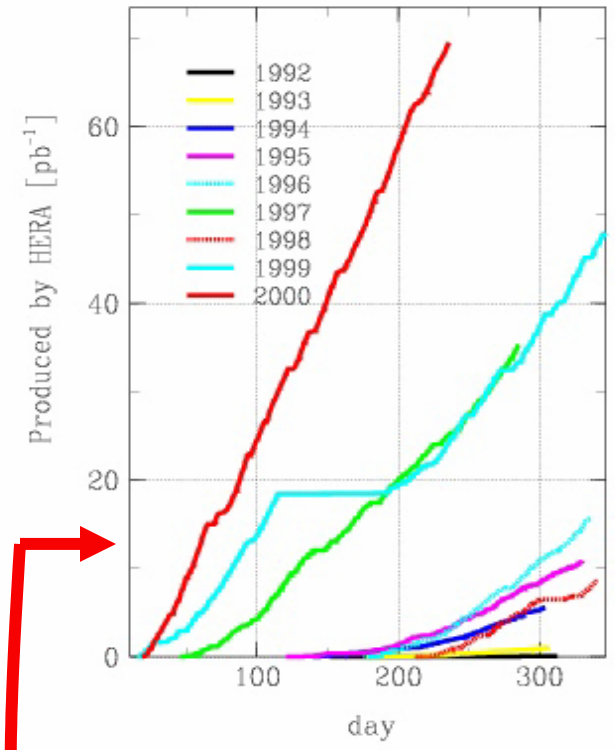
The Hera ep collider at Desy

Are quarks elementary

ep collisions allow to probe efficiently the structure of the quarks



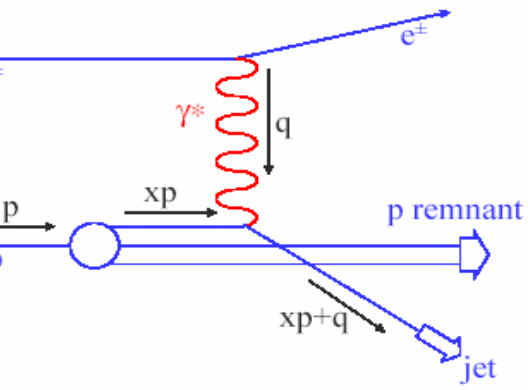
ing: 6 Km



'00 ~ 0.1 fb⁻¹ per experiment

'06 ~ 1 fb⁻¹ per experiment (?)

20 Feb 2004



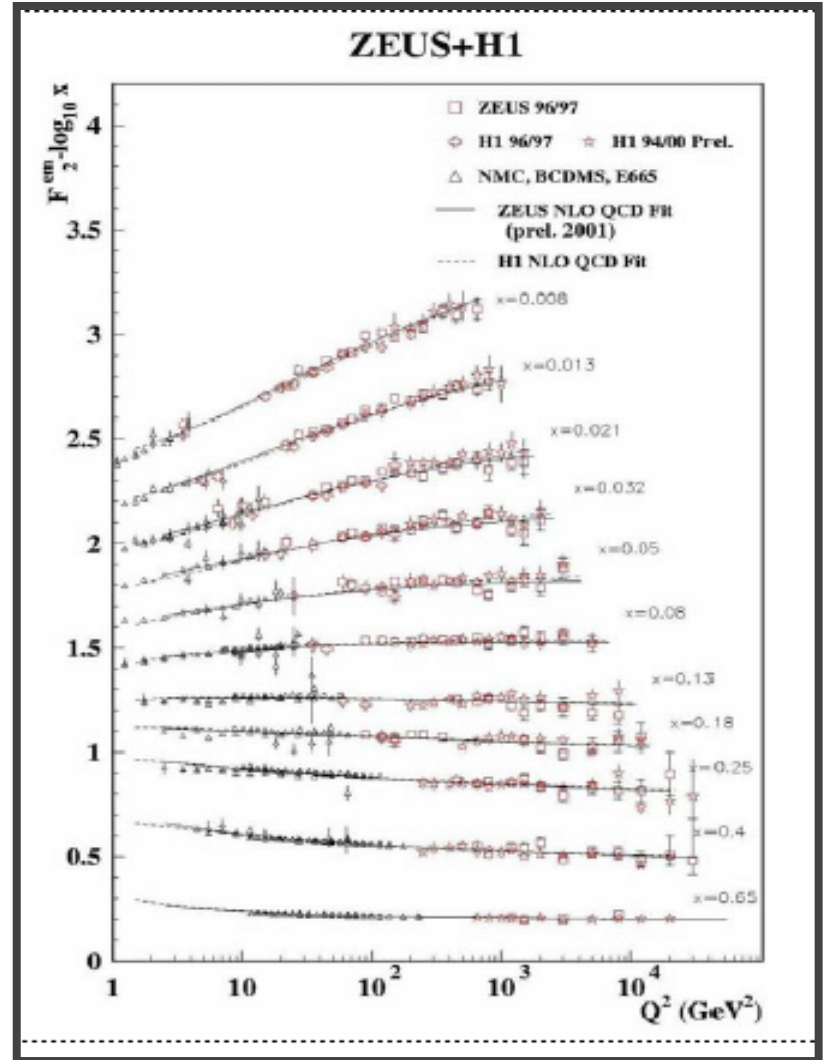
Describe the scattering in term of

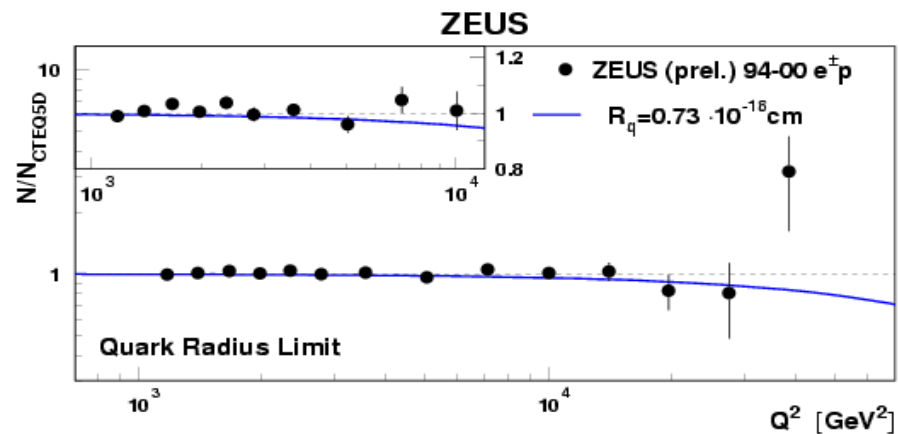
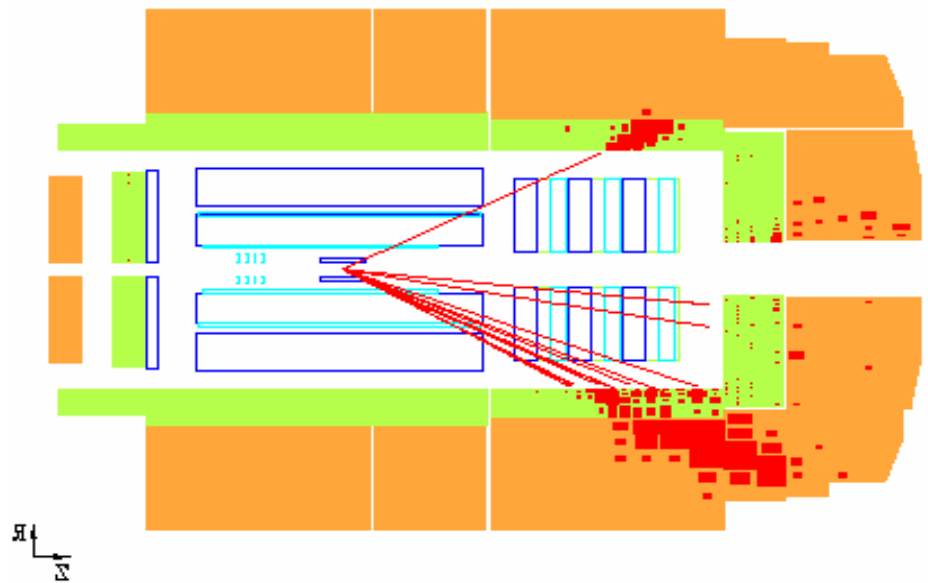
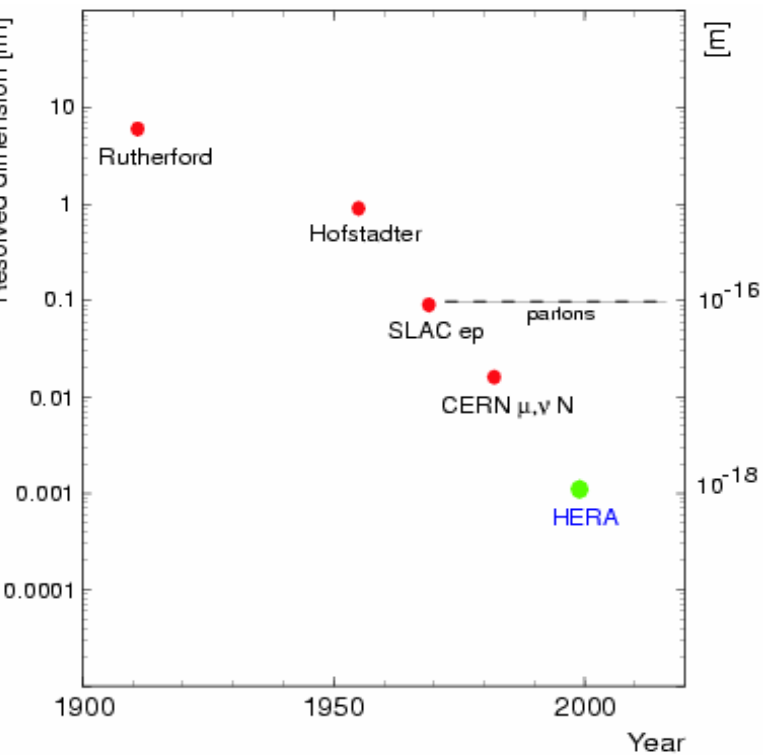
$$Q^2 = -q^2 \text{ and } x = \frac{Q^2}{2p \cdot q}$$

The cross section is expressed in term of the quark densities

$$\frac{d^2\sigma_{ep \rightarrow eX}}{dx dQ^2} \approx \frac{2\pi\alpha^2}{xQ^4} F_2(x, Q^2)$$

The accuracy of the measurement of angles and energies of leptons and jets is the challenge of the measurement to the cross section at high Q^2

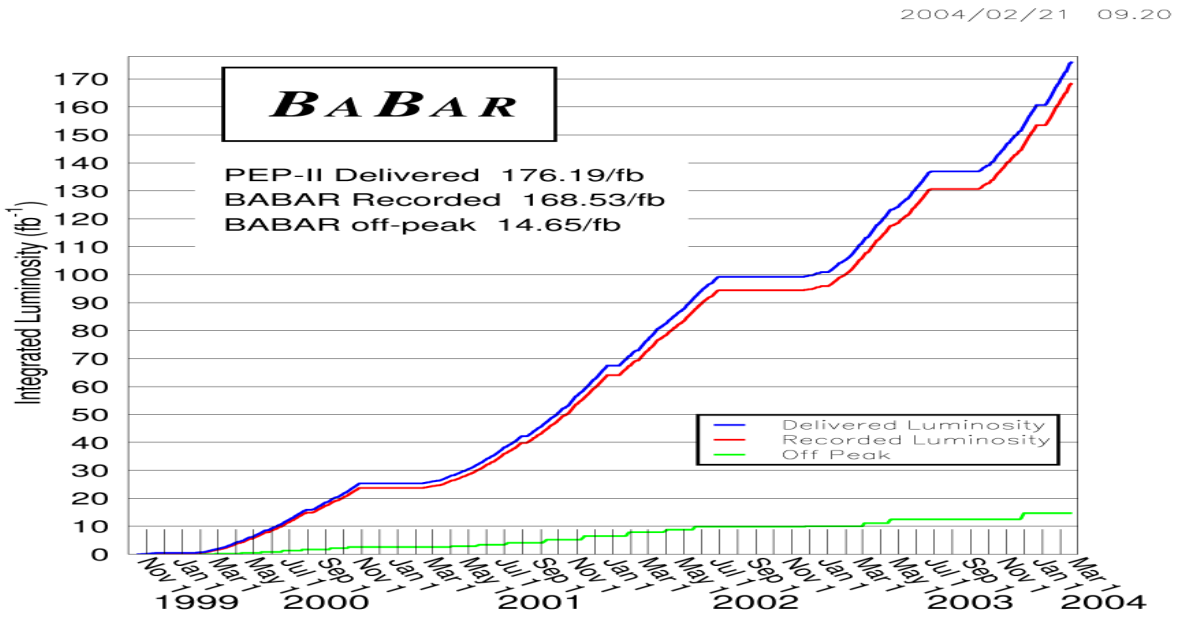




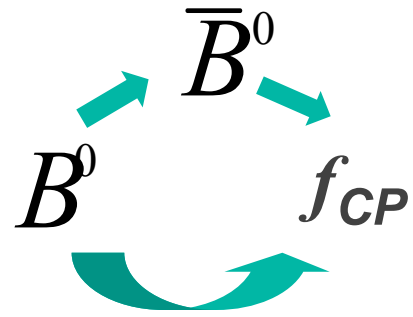
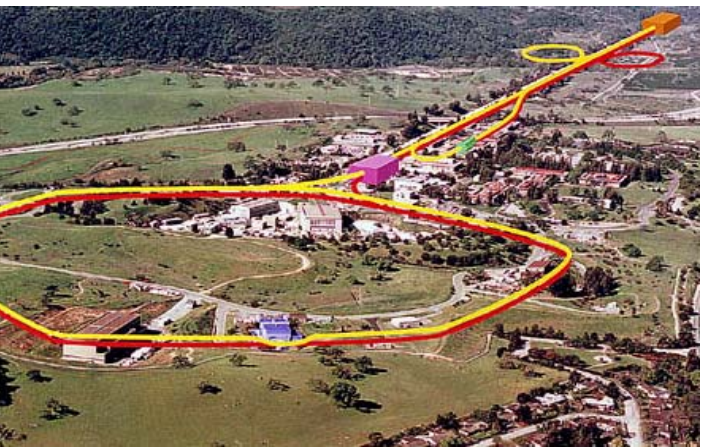
QCD with elementary quarks describes the scattering up to the highest accessible Q^2

The asymmetric B factories at Kek and Slac

Why three families?
Why matter?



~200.000.000 $B B_{\text{bar}}$ Events Collected



In the weak interaction u-type quarks couple to d-type quarks via the CKM matrix

CKM Matrix

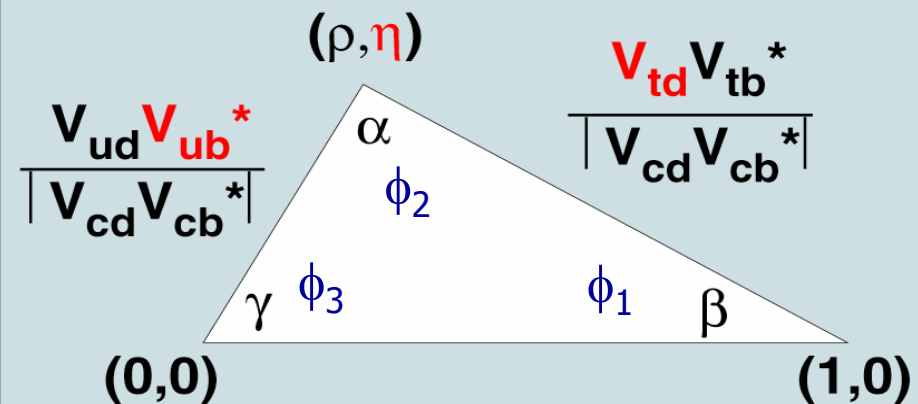
$$\begin{bmatrix} d' \\ s' \\ b' \end{bmatrix} = \begin{bmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{bmatrix} \begin{bmatrix} d \\ s \\ b \end{bmatrix}$$

$V^\dagger V = I$, and quark phases
 \Rightarrow 4 parameters

$$\begin{bmatrix} 1 - \frac{1}{2}\lambda^2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \frac{1}{2}\lambda^2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{bmatrix} + O(\lambda^4)$$

Unitarity Triangle

$$V_{ud}V_{ub}^* + V_{cd}V_{cb}^* + V_{td}V_{tb}^* = 0$$

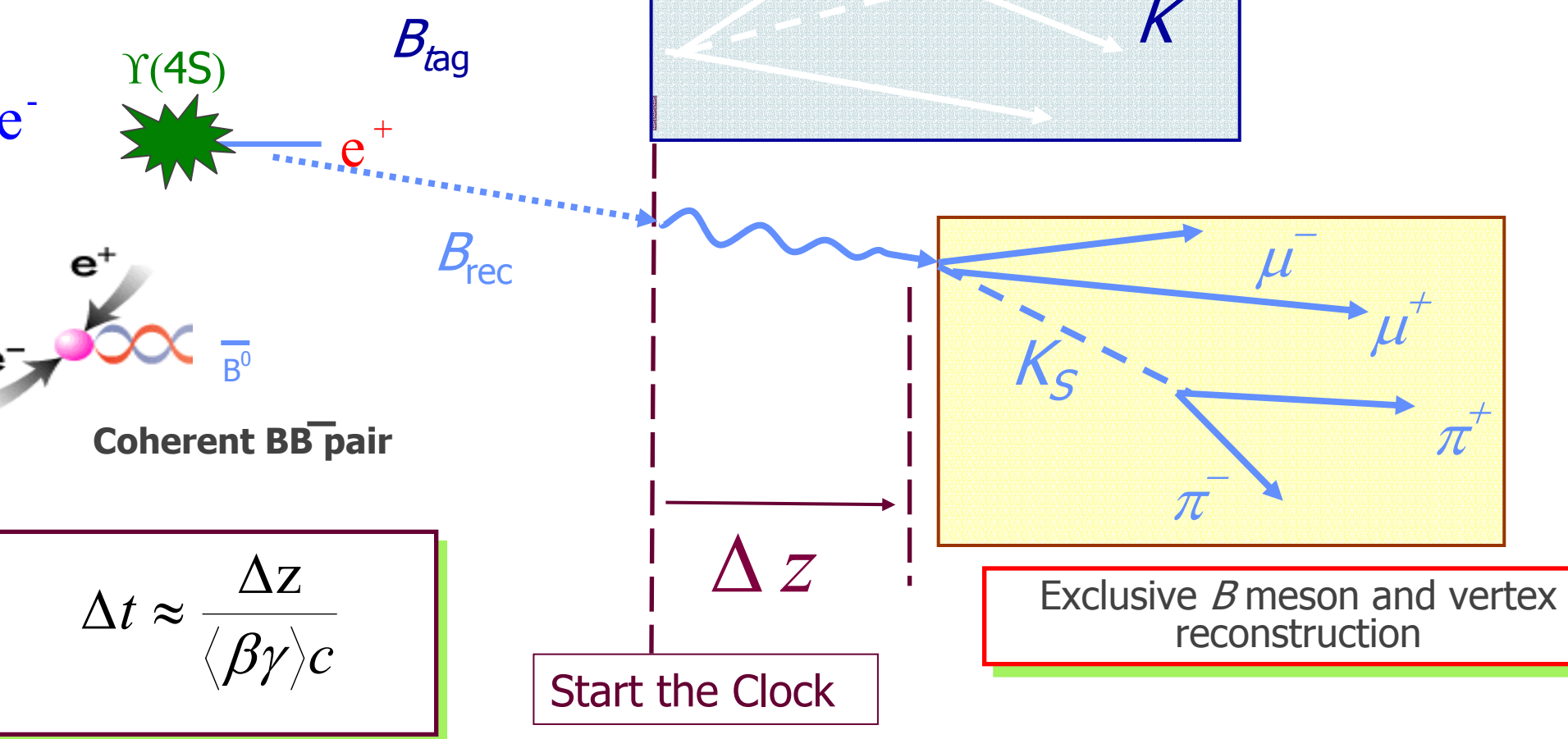


CP violation will arise from complex component of V_{ub} , V_{td}

Experimental technique at the $\Upsilon(4S)$ resonance

$$e^+ e^- \rightarrow \Upsilon(4S) \rightarrow B B^-$$

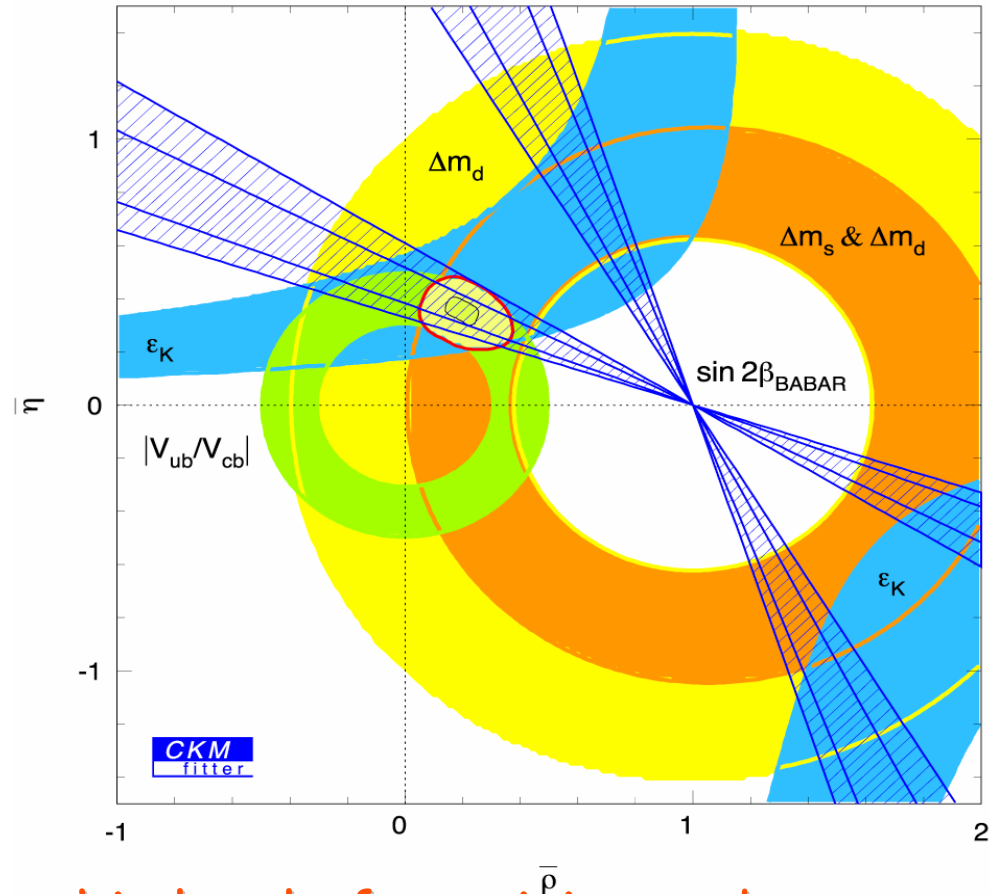
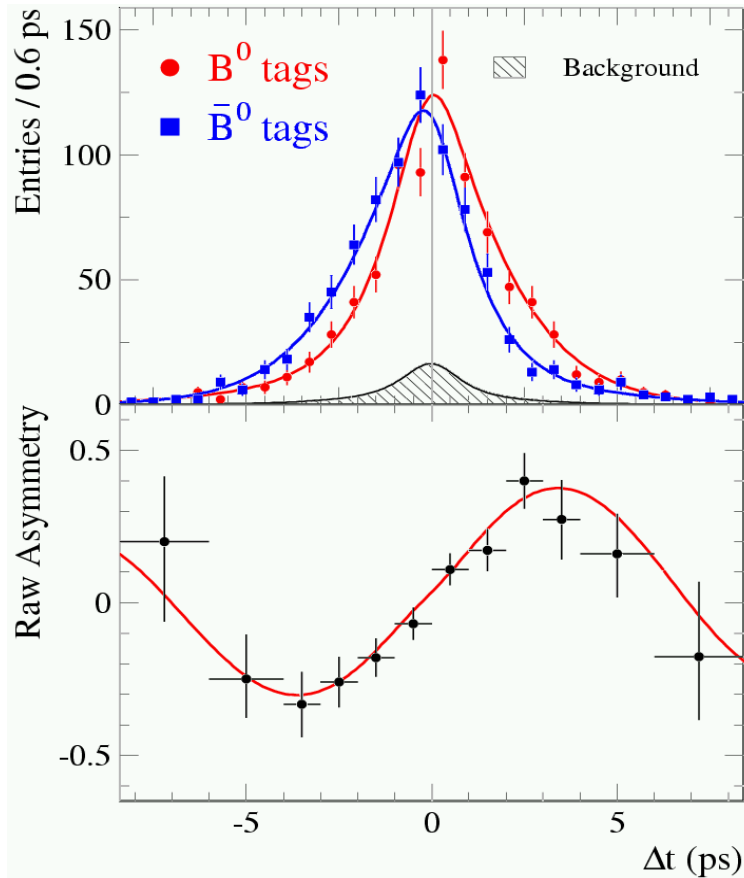
Boost: $\beta\gamma = 0.55$



$$\Delta t \approx \frac{\Delta z}{\langle \beta\gamma \rangle c}$$

Start the Clock

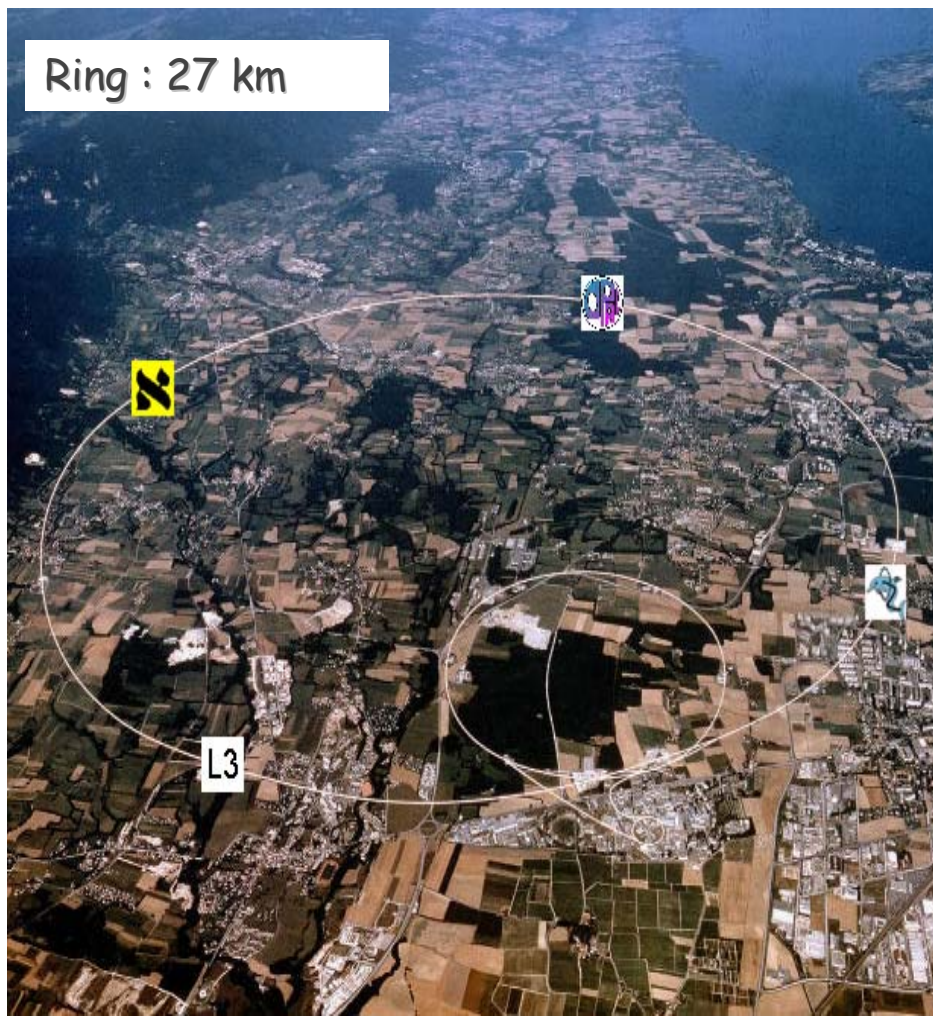
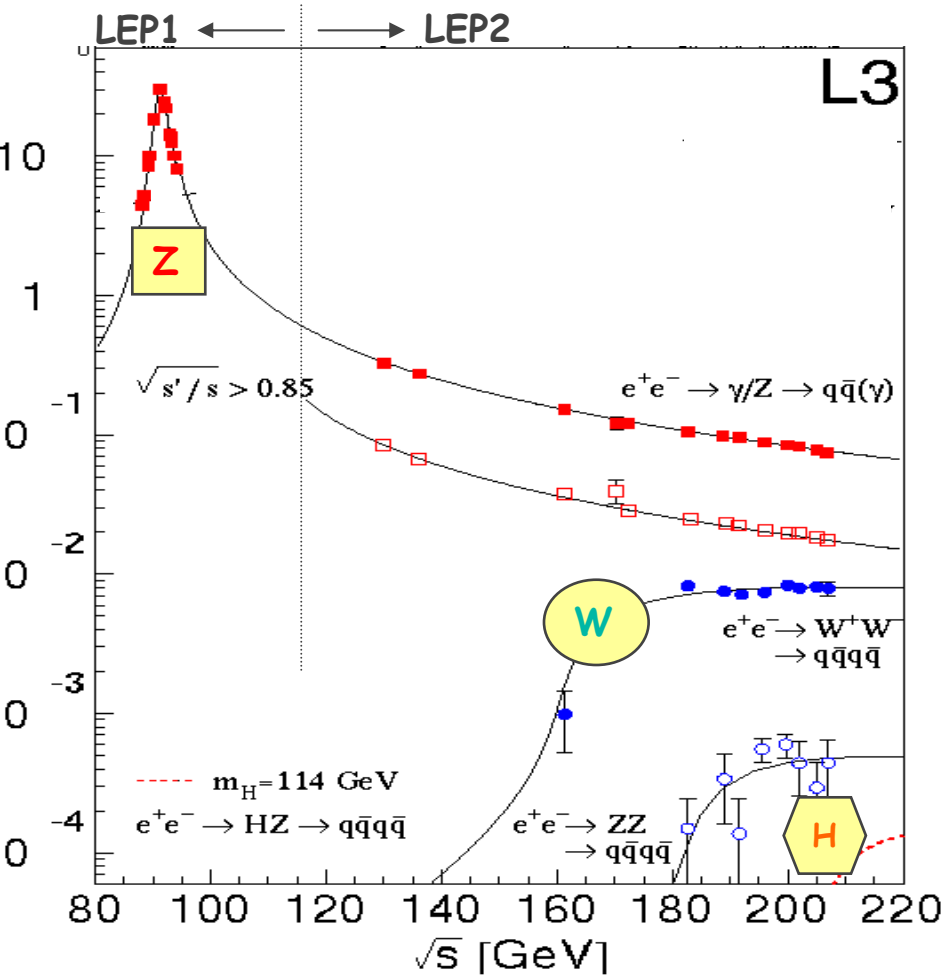
The measurement of the beta angle agrees at few percent level to its SM prediction based on other measured quantities



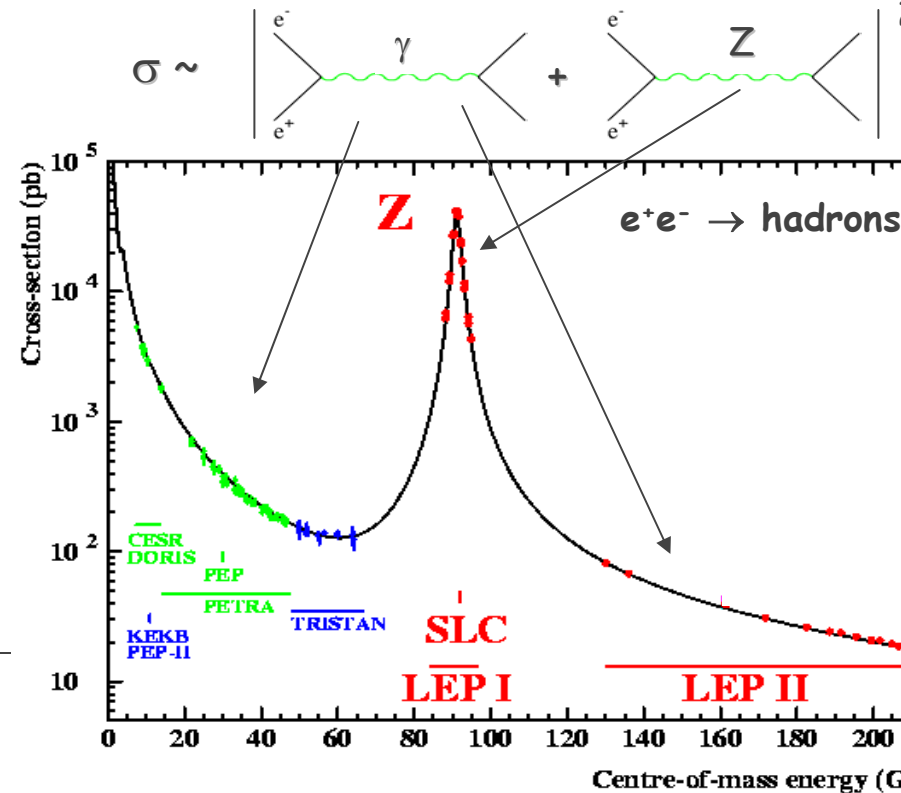
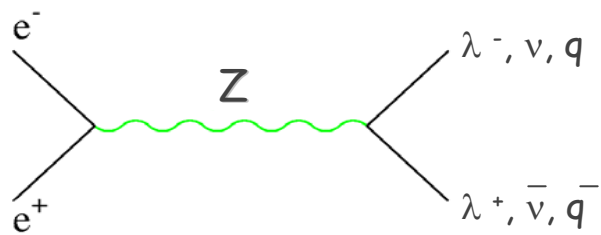
CKM matrix is unitary to this level of precision and incorporates CP violation with three generations

The LEP e+e- Collider at CERN

LEP1 ('89-'95) : $\sqrt{s} \approx m_Z \rightarrow 2 \cdot 10^7$ Z recorded \rightarrow precise Z measurements
 LEP2 ('96-2000) : $\sqrt{s} \rightarrow 209$ GeV \rightarrow WW production, m_W , search for Higgs and new particles



EP1 (CERN) and SLC (Stanford) e^+e^- Colliders start precision tests of SM at high energy $\sqrt{s} = E(e^-) + E(e^+) \approx m_Z \approx 90 \text{ GeV}$



achieved precision: better than 10^{-3}

measured observables:

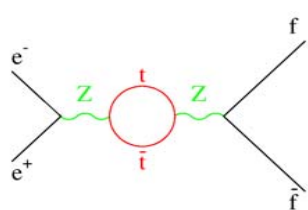
- m_Z, Γ_Z
- Z production cross-section
- all properties of Z couplings to fermions: e.g. decay modes, angular distributions etc..

WHY precision tests of the SM at high energy?

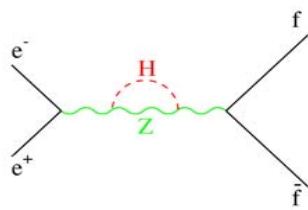
Best radiative quantum corrections (sensitive to heavy physics) :

lowest order

Examples of radiative corrections



$$\sim m_{\text{top}}^2$$



$$\sim \log m_H$$

$O_i \sim f_i (\alpha_{EM}, G_F, m_Z, m_{\text{top}}^2, \log m_H, \dots)$
 \rightarrow deduce masses of particles not directly produced

C. modify observables by $\approx \%$:

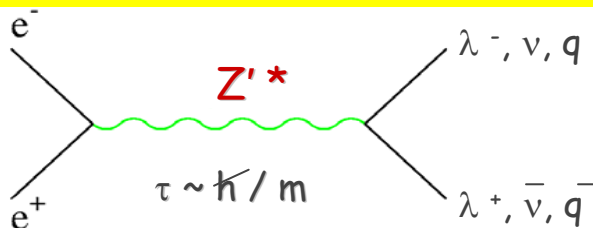
experimental precision of $\approx \%$ and improved theoretical needed

$m_{\text{top}} \sim 175 \text{ GeV}$ predicted by LEP/SLC in '94 before direct discovery at Tevatron pp Collider in '94-'95

New Physics can also contribute to loops (e.g. SUSY particles if light)

Beyond Z peak, search indirectly for New Physics by looking for deviations from SM

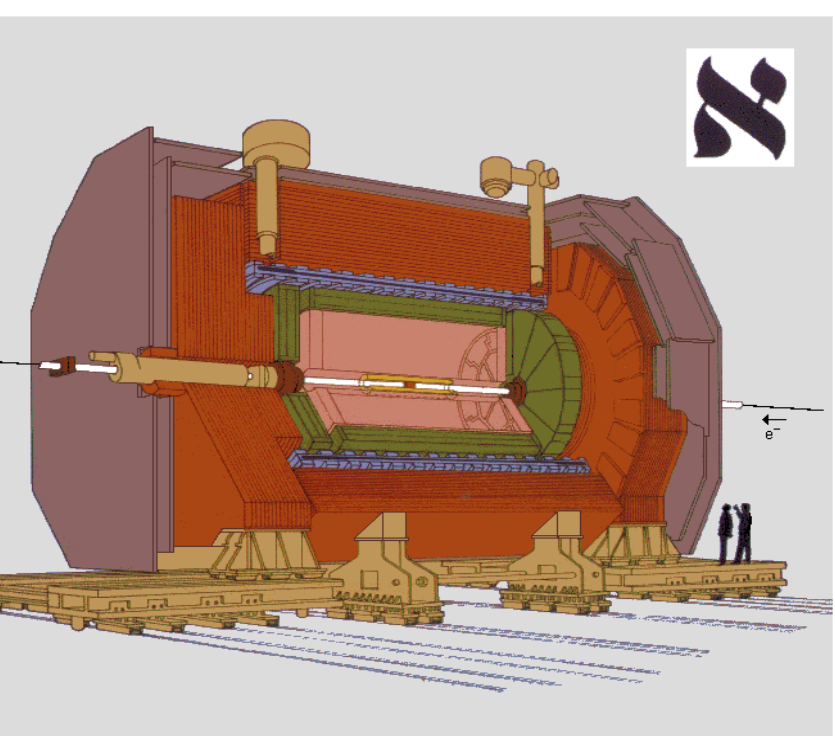
e.g. additional weak bosons



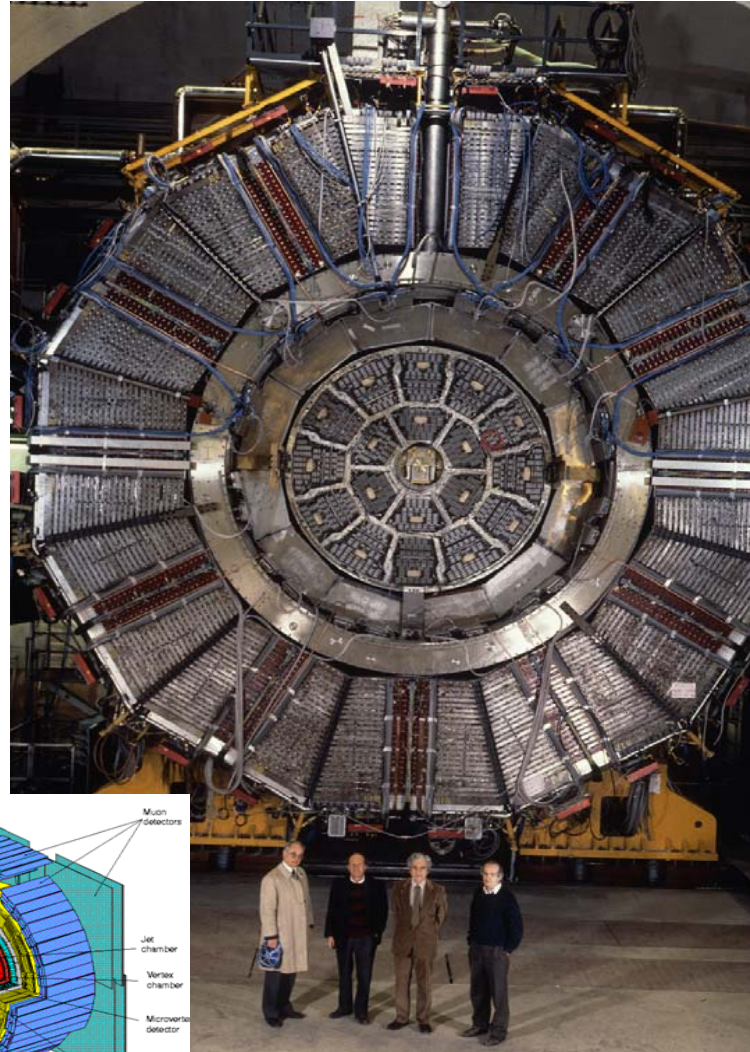
interferes with SM processes

\rightarrow deviations from SM expectations

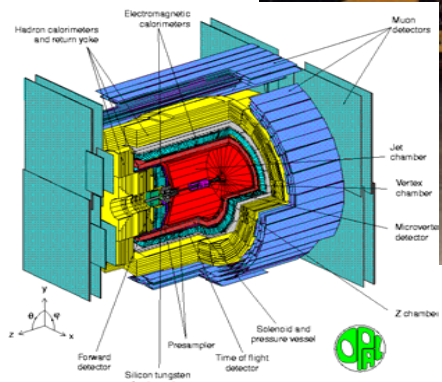
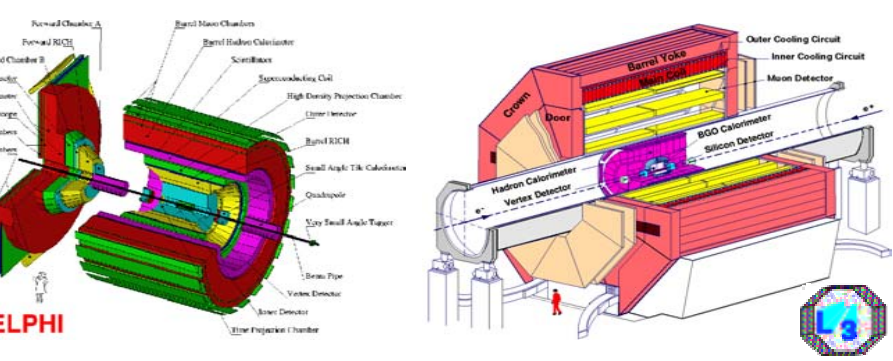
Z Lineshape: Final State Identification (I)



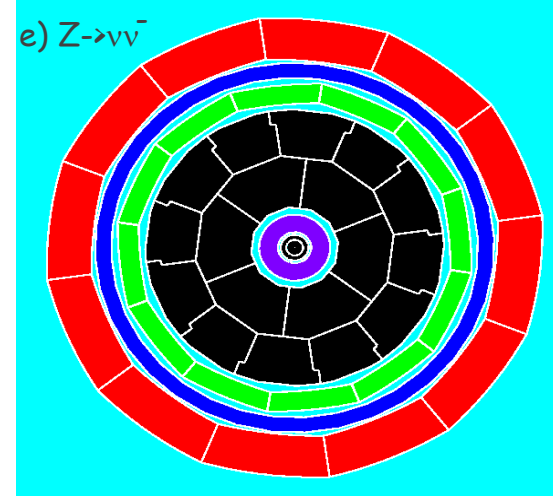
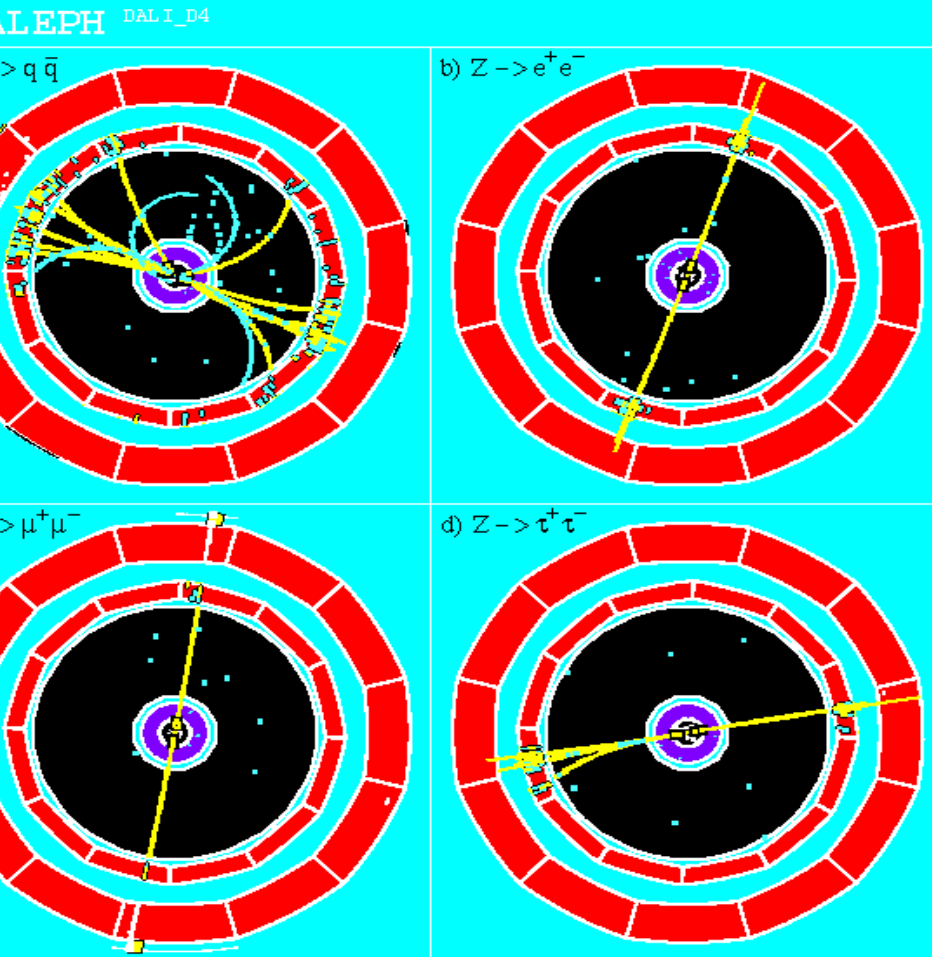
- Vertex Detector
- Inner Tracking Chamber
- Time Projection Chamber
- Electromagnetic Calorimeter
- Superconducting Magnet Coil
- Hadron Calorimeter
- Muon Chambers
- Luminosity Monitors



The ALEPH Detector



Z Lineshape: Final State Identification (II)



• $Z \rightarrow \nu\bar{\nu}$:
Not detectable.

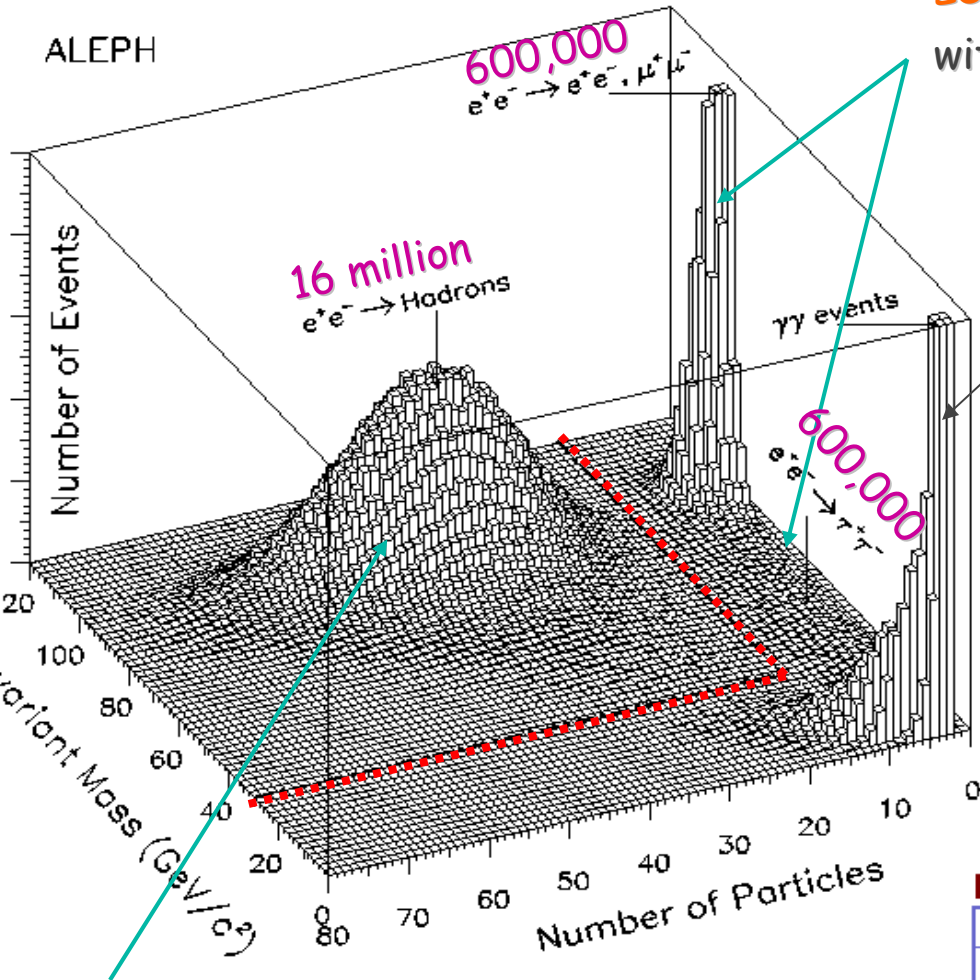
• $Z \rightarrow \tau^+\tau^-$: Two low multiplicity jets + missing energy carried by the decay neutrinos

$Z \rightarrow q\bar{q}$: Two jets, large particle multiplicity.

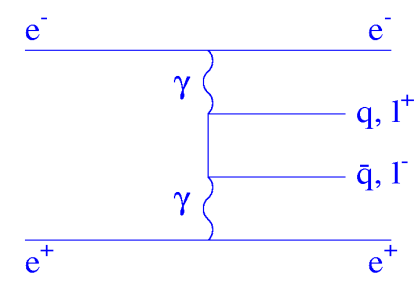
$Z \rightarrow e^+e^-, \mu^+\mu^-$: Two charged particles (e or μ .)

Channel	Partial Width	Branching Ratio
Hadrons	1.739 GeV	70%
Neutrinos	0.497 GeV	20%
Leptons	0.250 GeV	10%

Z Lineshape: Final State Identification (III)



Leptonic decays: Low multiplicity, with (τ) or without (e, μ) missing energy



$\gamma\gamma$ Collisions:
Low multiplicity,
Low mass

- Selections with
- High Efficiency;
 - High Purity;

Count events : Easy?

Hadronic decays:
High multiplicity
High mass

Note: Also need precise Luminosity determination

FLORIDA VOTE COUNT TOTALS

PRESIDENT	Nov. 7	First Recount	Certified
R Bush	2,909,176	2,911,872	2,912,790
D Gore	2,907,451	2,910,942	2,912,253
Bush Lead	1725	930	537

Source: State of Florida. Systematic Uncertainty ~ 0.1%
25 electoral votes at stake

Z Lineshape: Results (III)

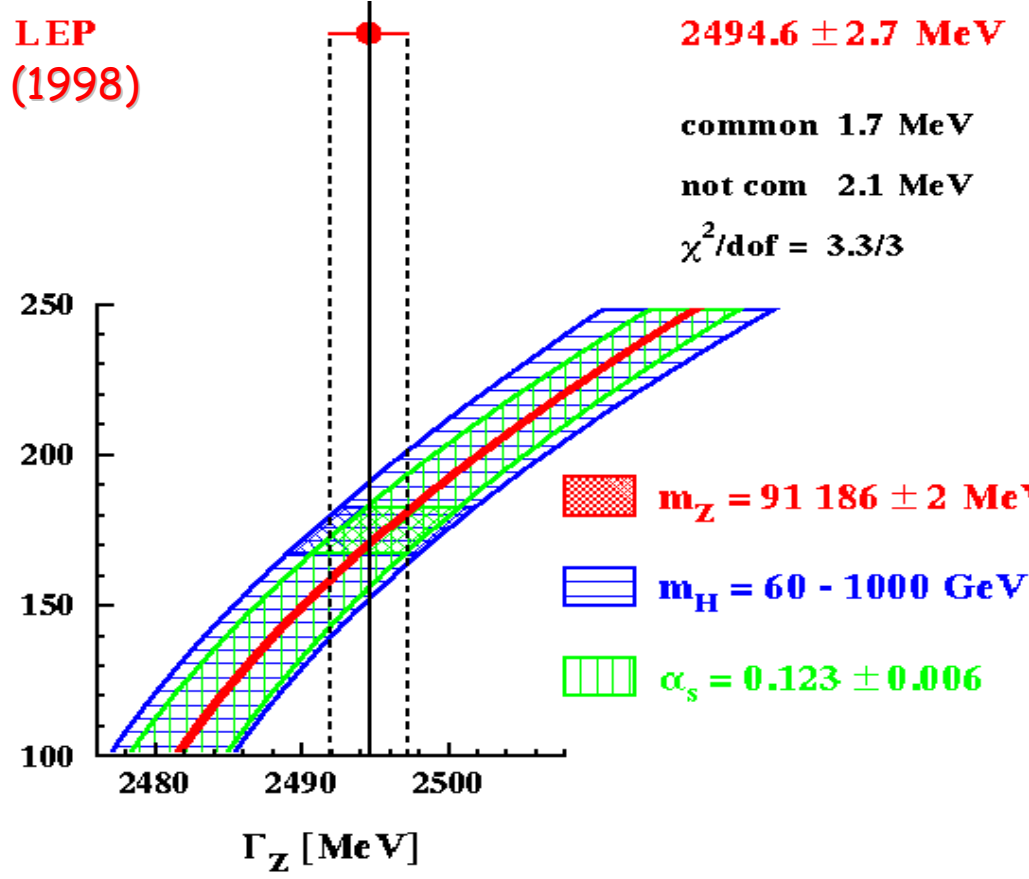
Obs.	Value	Error
m_Z	91187.5	2.1 MeV
Γ_Z	2495.2	2.3 MeV
σ^0	41.540	0.037 nb
R_l	20.767	0.025

10⁻³

$$\sigma_{\text{had}}^0 / \sigma_l^0$$

500 MeV
in 1989

$$\Gamma_Z \propto (1 + \Delta\rho)^2$$



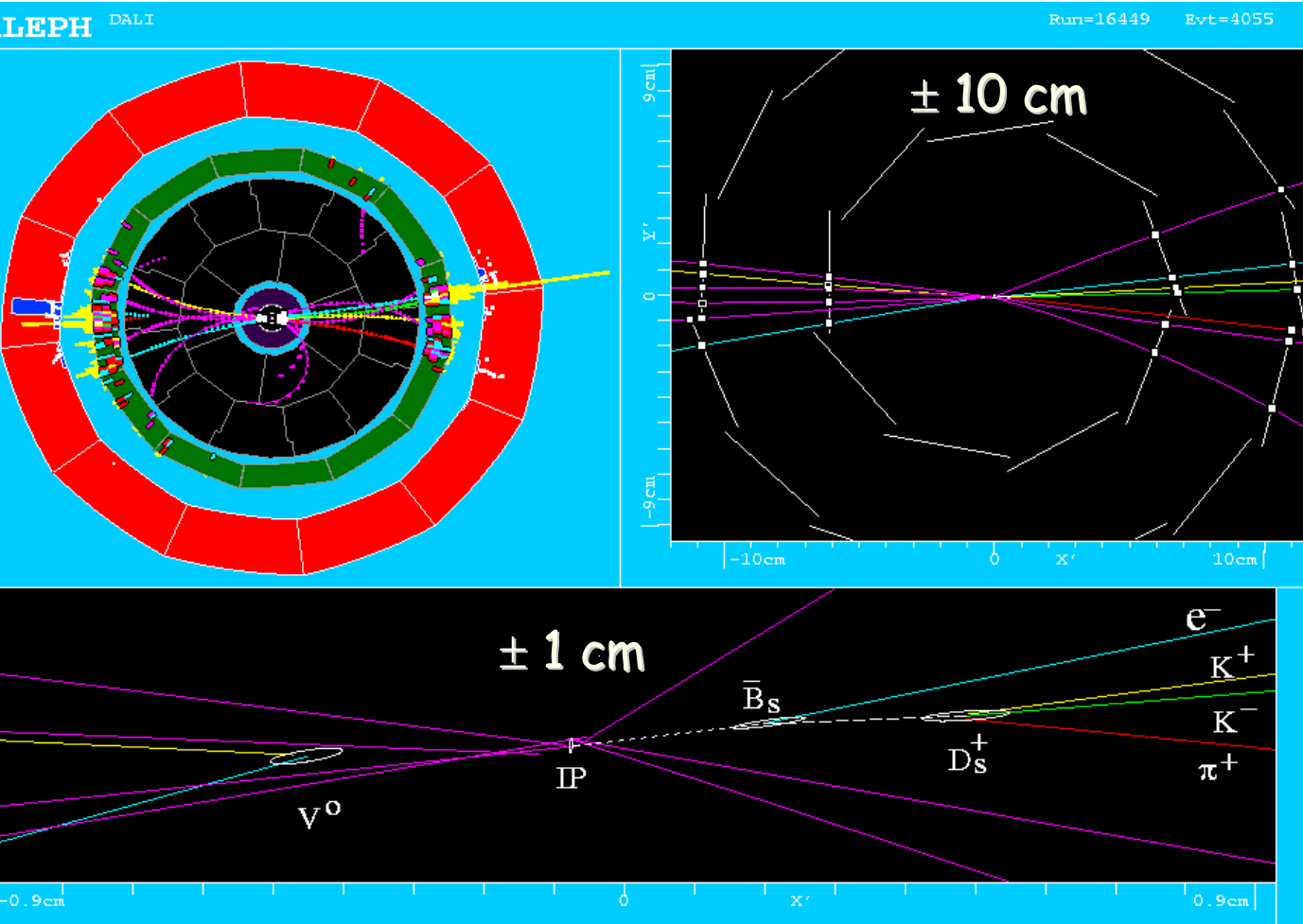
With this measurement alone:

$m_{\text{top}} \sim 165 \pm 25 \text{ GeV}/c^2$

(+small sensitivity to m_H)

Heavy Flavour Rates: Identification (I)

b- and c-hadrons decay weakly towards c- and s-hadrons, with a macroscopic lifetime (1.6 ps for b's), corresponding to few mm's at LEP



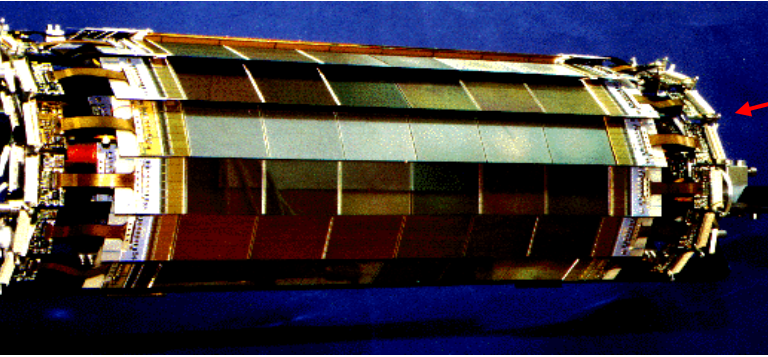
3d-vertexing determines secondary and tertiary vertices.

High resolution is crucial

Impact parameters of reconstructed tracks allow b quarks to be tagged with very good purity.

Mass of secondary vertex tracks is a very powerful discriminator of flavour (b, c, and light quarks):
 $m_b \sim 5 \text{ GeV}/c^2$, and
 $m_c \sim 1.5 \text{ GeV}/c^2$.

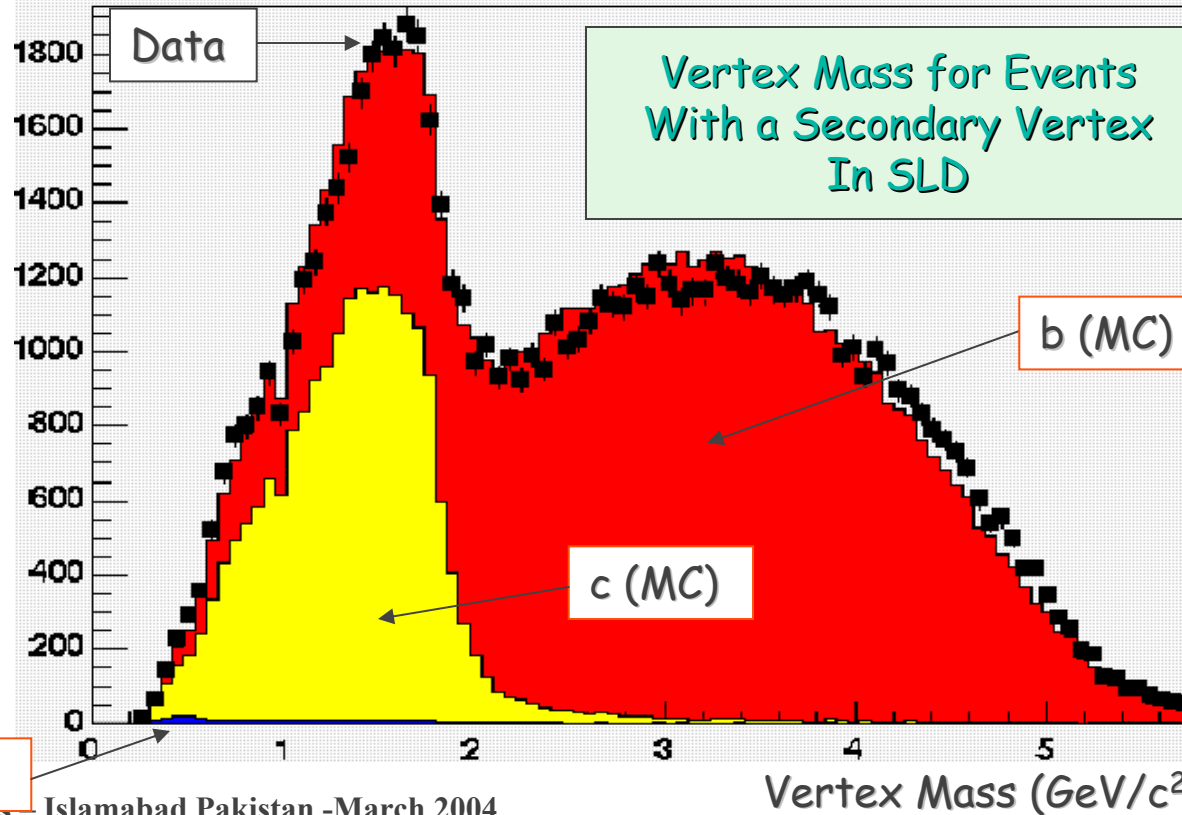
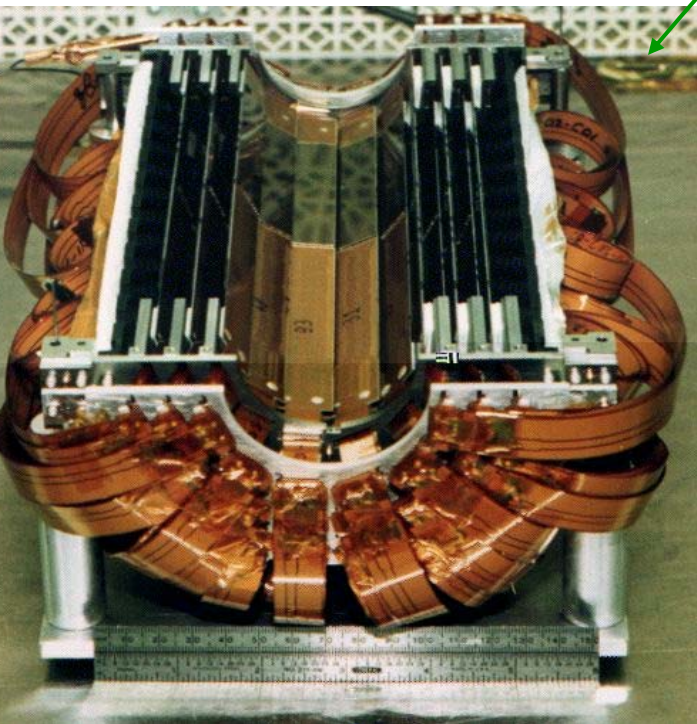
Heavy Flavour Rates: Identification (II)



Vertex detectors (Si μ -strips, CCD's, pixels):

- At LEP: inner radius 6 cm, good resolution;
- At SLC: inner radius 2.3 cm, superior resolution.

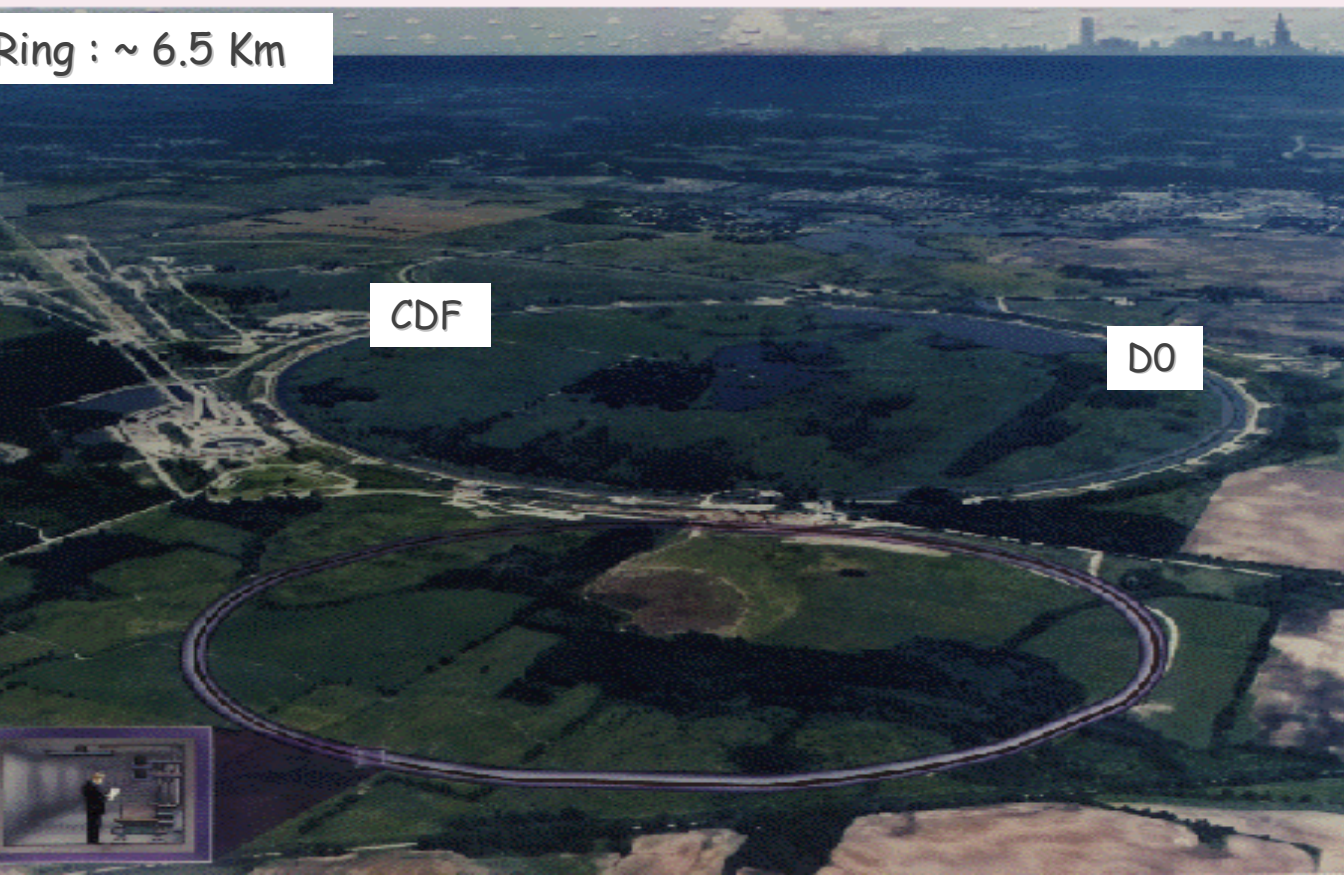
SLD can do both b- and c-tagging with good purity.



The Tevatron $p\bar{p}$ Collider at Fermilab

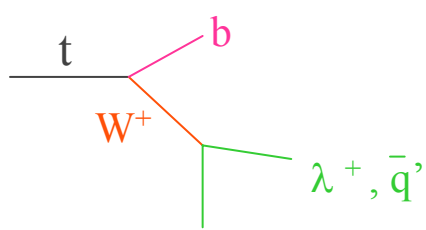
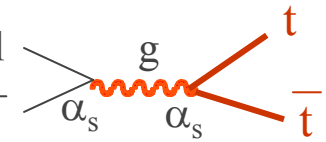
$\sqrt{s} \approx 2 \text{ TeV}$

Ring : $\sim 6.5 \text{ Km}$



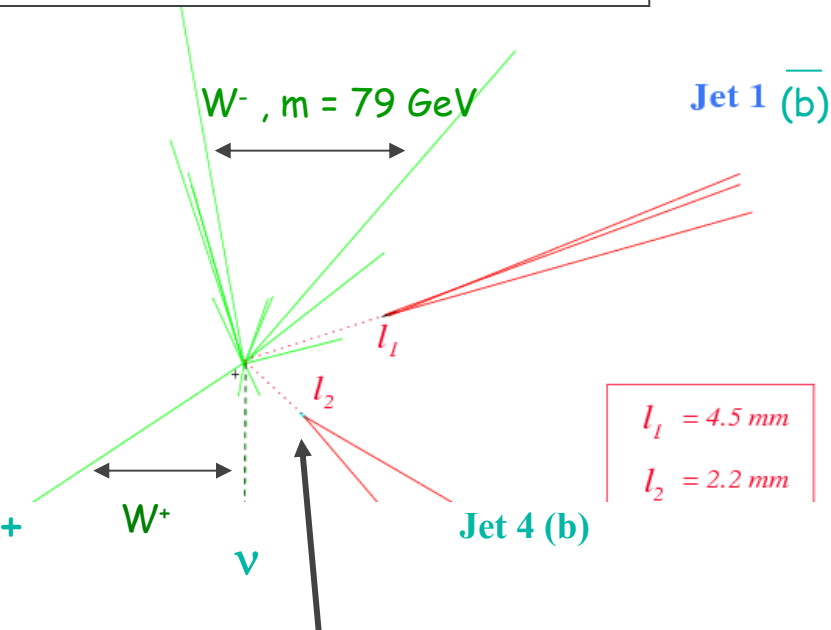
Run 1 ('89-'96) : ≈ 200 top events \rightarrow discovery of top
 $\approx 80\,000$ W events measurement of m_W and m_{top}
Run 2 ('01-'07?) ≥ 100 times more data \rightarrow }
better measurements of m_W and m_{top} ,
searches for Higgs and new particles

The top quark at the Tevatron



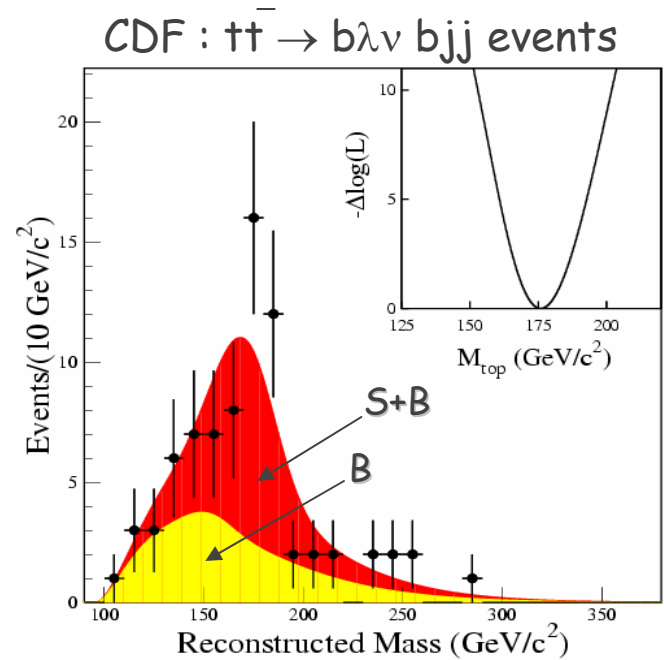
Heaviest particle observed so far
(and $m_{\text{top}} - m_b \sim 170 \text{ GeV}$) \rightarrow clues
about origin of masses?

$t\bar{t} \rightarrow bW \bar{b}W \rightarrow b\lambda\nu bjj$ event
from CDF data



Secondary vertices

(b-hadrons) $\sim 1.5 \text{ ps} \rightarrow$ decay at
few mm from primary vertex
detected with high-granularity
Si detector (b-tagging)

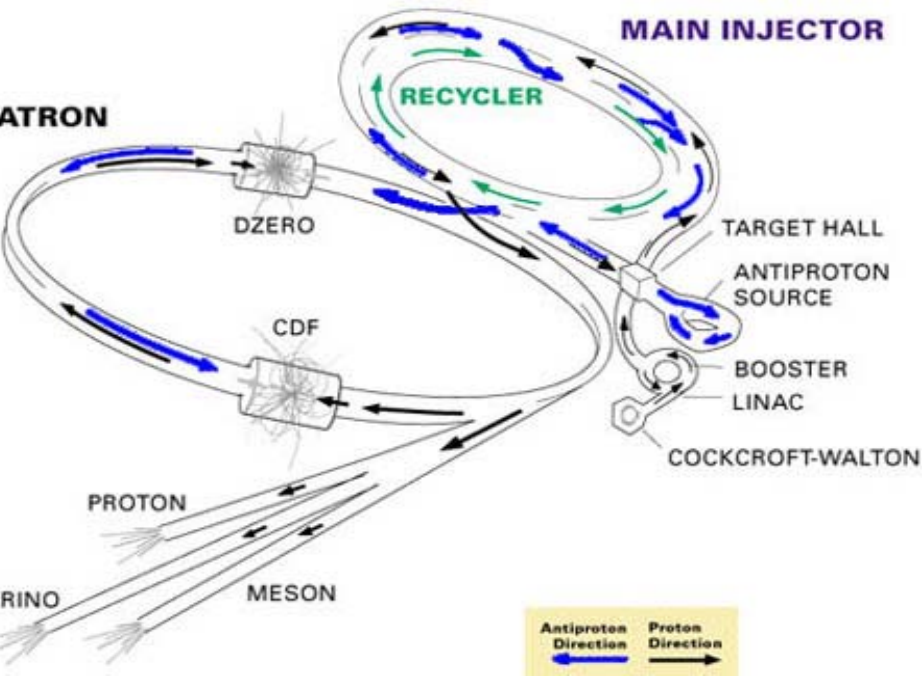


$m_{\text{top}} (\text{CDF} + \text{D0}) = 174.3 \pm 5.1 \text{ GeV}$

3%

statistics, calorimeter
calibration

Upgraded Tevatron

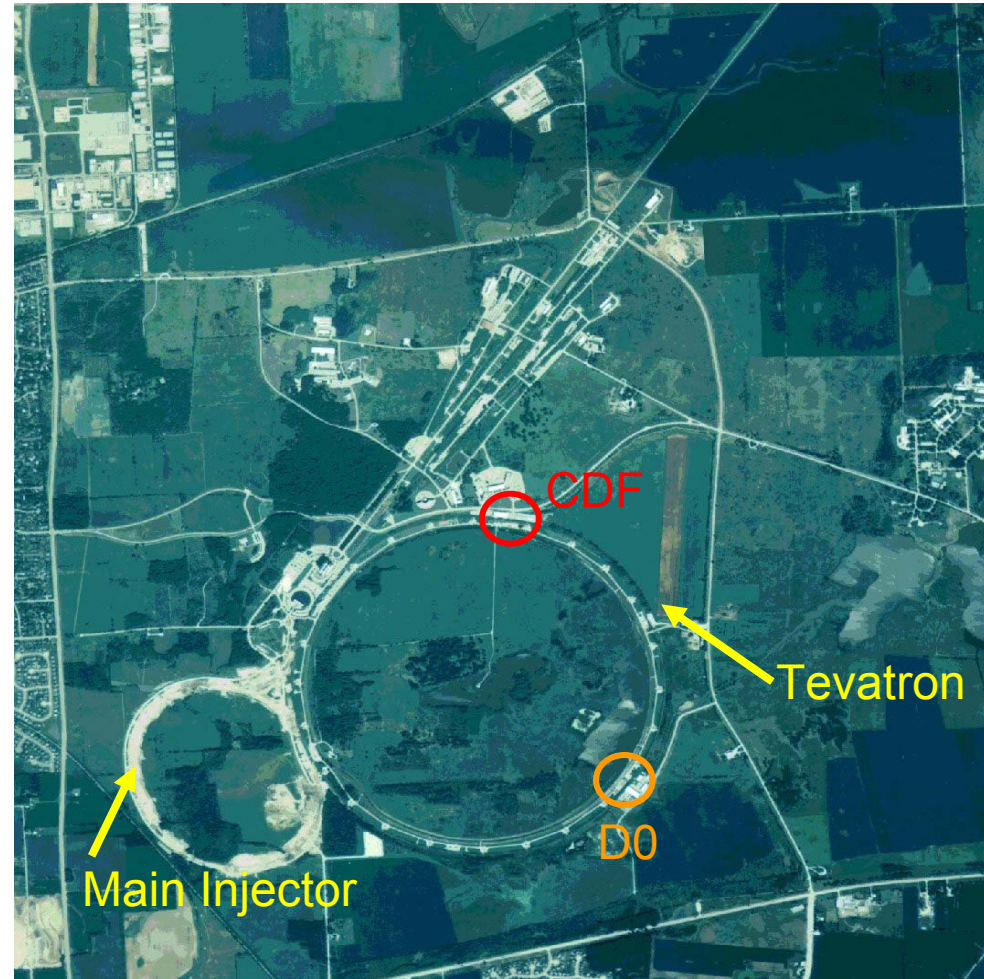


New Main Injector:

- Improve p-bar production

Recycler ring:

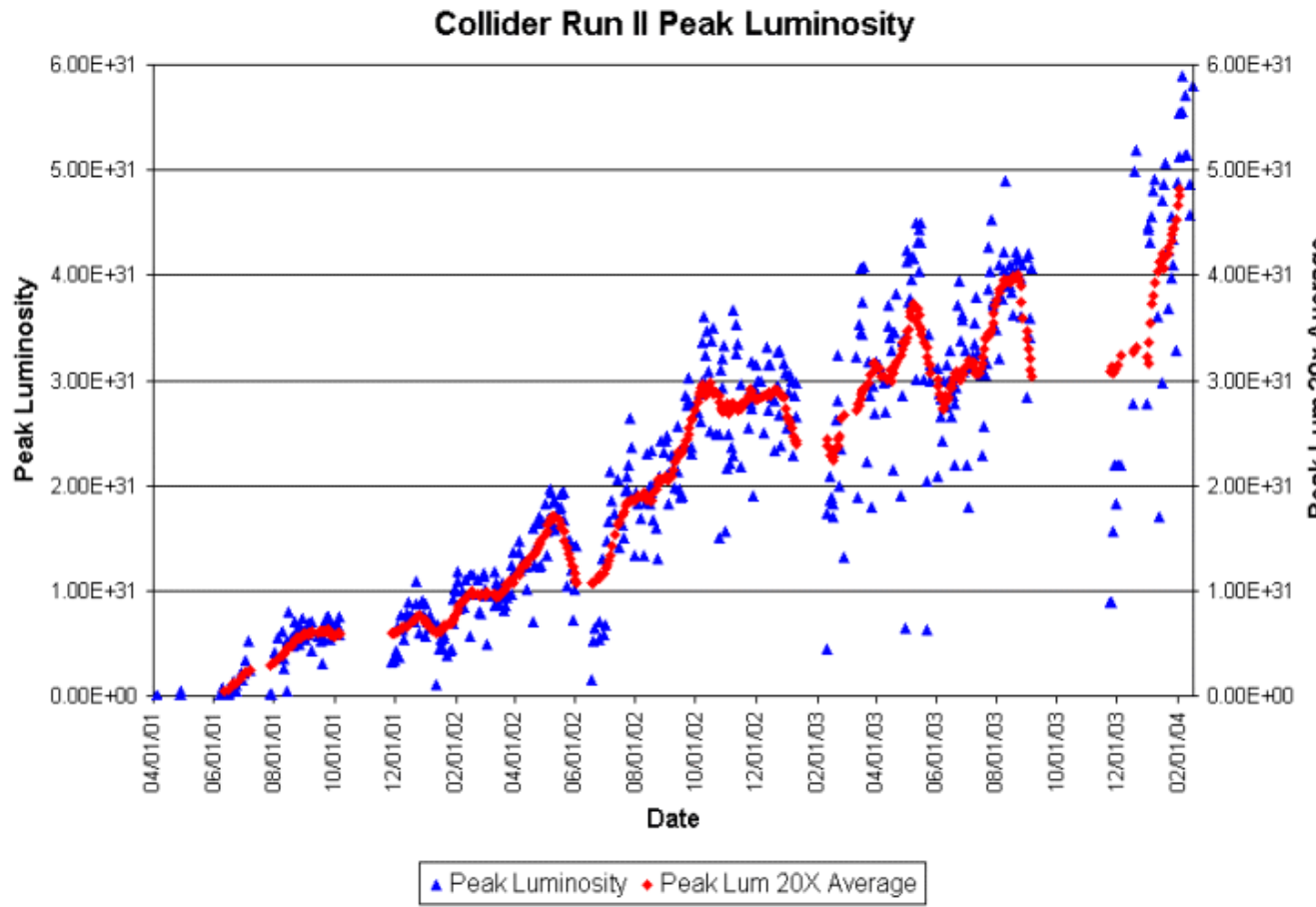
- Reuse p-bars! (current crisis)



TeV Luminosity (current situation)

Peak luminosity still low
but improving

– Best 6×10^{31}

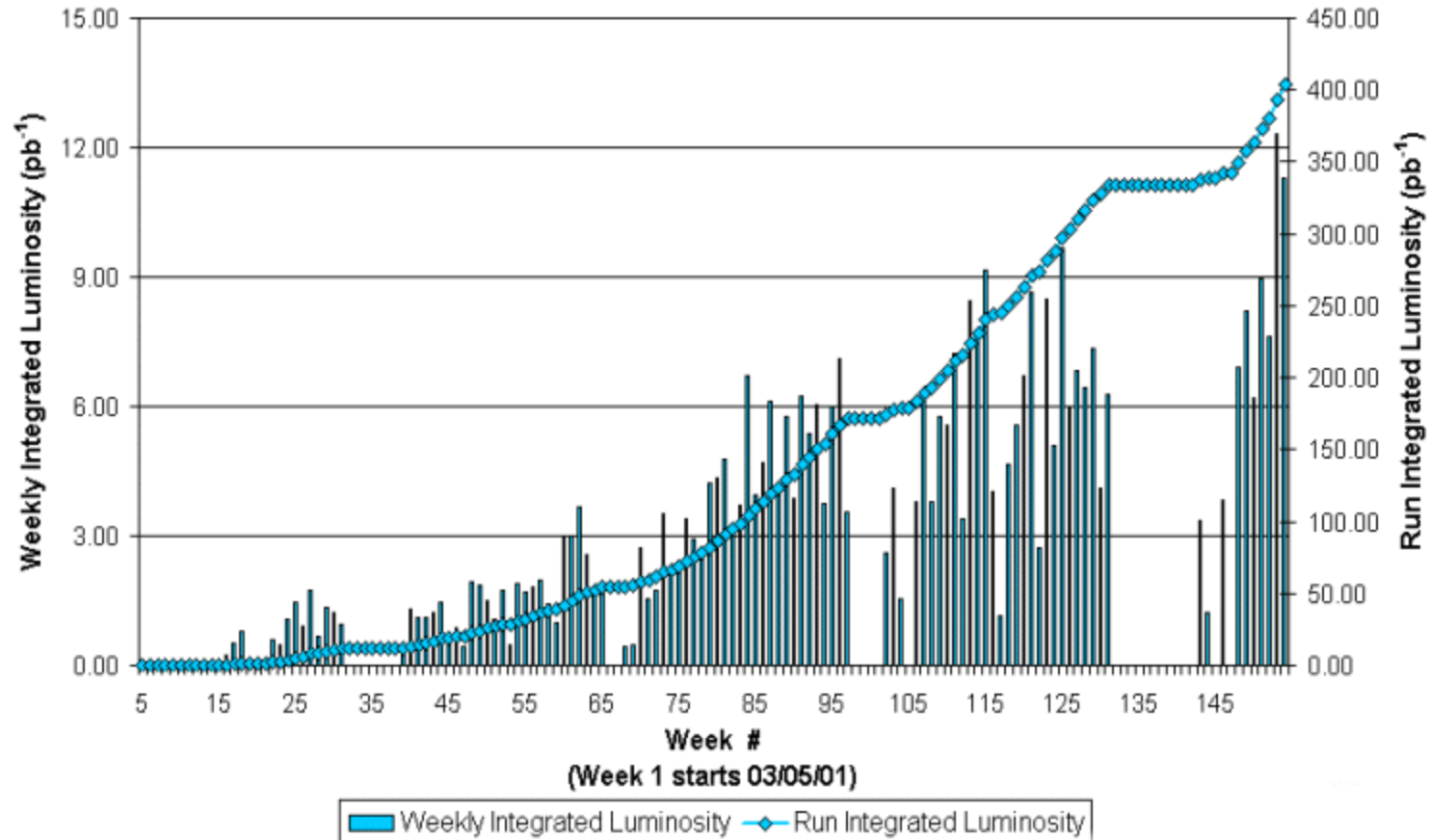


Tev luminosity

6 pb⁻¹ /wk typical

– Peak 12 pb⁻¹ /wk

Collider Run II Integrated Luminosity



Trigger

Some basic numbers at Tevatron

- Bunch crossing frequency = 2.5 MHz
- > 1 interaction/crossing
- Max data logging frequency ~ 100 Hz

Need ~ 10^4 rejection factor

provided by **trigger system**

- Custom electronics on detector (L1)
- ~1 room full of custom electronics (L1/2)
- ~1 room full of dedicated PC's (L3)

Significant fraction of analysis done at trigger level

